

# MA542: Differential Equations

## Lecture - 38

18/04/2022

# Finite Vibrating String with no External Force

- Recall the finite string problem in a computational domain  $(x, t) \in [0, L] \times [0, \infty)$

► The governing equation:

$$u_{tt} = c^2 u_{xx}, \quad (x, t) \in (0, L) \times (0, \infty). \quad (1)$$

► The boundary conditions for all  $t > 0$ :

$$u(0, t) = 0, \quad u(L, t) = 0. \quad (2)$$

► The initial conditions for  $0 \leq x \leq L$ :

$$u(x, 0) = \phi(x), \quad u_t(x, 0) = \psi(x). \quad (3)$$

# Formal Solution of the Finite Vibrating String Problem

The solution is given by

$$u(x, t) = \sum_{n=1}^{\infty} \sin \frac{n\pi x}{L} \left[ A_n \cos \frac{n\pi ct}{L} + B_n \sin \frac{n\pi ct}{L} \right], \quad (4)$$

with

$$\begin{aligned} A_n &= \frac{2}{L} \int_0^L \phi(x) \sin \frac{n\pi x}{L} dx, \quad n = 1, 2, 3, \dots, \\ B_n &= \frac{2}{n\pi c} \int_0^L \psi(x) \sin \frac{n\pi x}{L} dx, \quad n = 1, 2, 3, \dots \end{aligned}$$

## IBVP for Vibrating string (Contd.)

The individual displacement for each  $n$  in (4) is referred as the  $n$ -th eigenfunction or the  $n$ -th normal mode of the vibrating string.

The  $n$ -th normal mode vibrates with a period of  $\frac{2L}{nc}$  seconds which corresponds to a frequency of  $\frac{nc}{2L}$  cycles per second.

Since  $c^2 = T/\rho$ , where  $T$  is the tension and  $\rho$  is the density of the string, the frequency is

$$\frac{n}{2L}(T/\rho)^{1/2}.$$

Hence, if a string on a musical instrument is vibrating in a normal mode, its pitch may be sharpened (frequency increased) by either decreasing the length  $L$  of the string or increasing the tension in the string.

## IBVP for Vibrating string (Contd.)

The first normal mode  $n = 1$  vibrates with the lowest frequency

$$\frac{1}{2L}(T/\rho)^{1/2}.$$

This is called the *fundamental frequency* of the string.

If the string is made to vibrate in a higher mode, the frequency is increased by an integer multiple and this corresponds to the production of a musical harmonic or overtone.

When a vibrating system has multiples of fundamental frequency, say in a violin, then music is produced.

When a vibrating system has frequencies which are not integer multiples of fundamental frequency, then noise is produced.

# IBVP for Vibrating string (Contd.)

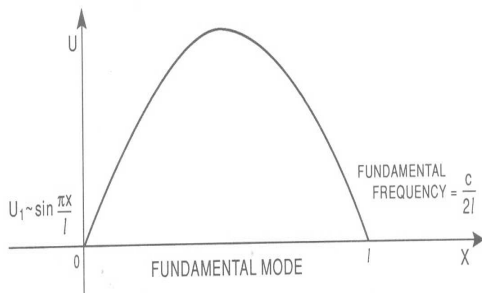


Figure : Fundamental mode of a vibrating string

# IBVP for Vibrating string with gravity(Contd.)

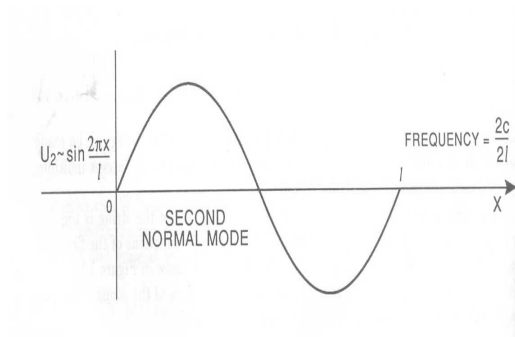


Figure : Second normal mode of a vibrating string

## IBVP for Vibrating string (Contd.)

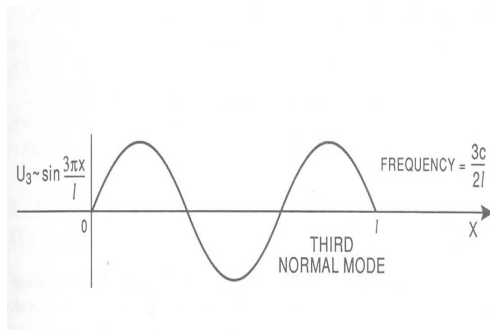


Figure : Third normal mode of a vibrating string



## IBVP for Vibrating string (Contd.)

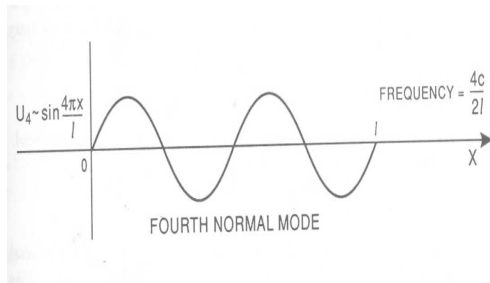
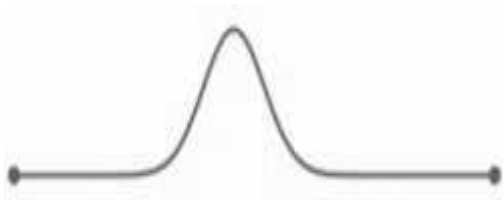


Figure : Fourth normal mode of a vibrating string

## IBVP for Vibrating string (Contd.)



**Figure :** A pulse traveling through a string with fixed endpoints as modeled by the wave equation

# Solution of the Finite Vibrating String Problem: Example

**Example:** For a string of length  $L$  stretched between the points  $x = 0$  and  $x = L$ , find the vibration in the string subject to the following initial conditions:

$$u(x, 0) = \sin(\pi x/L) + 1/2 \sin(3\pi x/L), \quad u_t(x, 0) = 0.$$

**Solution:** Here, initial conditions are

$$\phi(x) = \sin(\pi x/L) + 1/2 \sin(3\pi x/L), \quad \psi(x) = 0.$$

Therefore,  $B_n = 0$ .

and

$$A_n = \frac{2}{L} \int_0^L \left( \sin \frac{\pi x}{L} + \frac{1}{2} \sin \frac{3\pi x}{L} \right) \sin \frac{n\pi x}{L} dx.$$

Due to the orthogonality of the set  $\{\sin \frac{n\pi x}{L} : n = 1, 2, \dots\}$ , only  $A_1$  and  $A_3$  are non-zero, and they are found as  $A_1 = 1, A_3 = 1/2$ .

Therefore, the solution of the IBVP is

$$u(x, t) = \sin \frac{\pi x}{L} \cos \frac{\pi ct}{L} + \frac{1}{2} \sin \frac{3\pi x}{L} \cos \frac{3\pi ct}{L}.$$

## Finite Vibrating String with Gravity

We consider an external force due to the gravitational acceleration  $g$  only (consider a string oriented horizontally). Then the one-dimensional wave equation becomes

$$u_{tt} = c^2 u_{xx} - g, \quad 0 < x < L, \quad t > 0. \quad (5)$$

We seek to find

the displacement of the string at any position  $x$  and at any time  $t$  subject to the following boundary condition (for  $t > 0$ ) and initial conditions ( $0 \leq x \leq L$ ):

$$u(0, t) = 0, \quad (6a)$$

$$u(L, t) = 0, \quad (6b)$$

and

$$u(x, 0) = \phi(x), \quad (7a)$$

$$u_t(x, 0) = \psi(x). \quad (7b)$$

## Finite Vibrating String with Gravity (Contd.)

Due to the presence of the term  $g$  in equation (5), the equation has now become non-homogeneous and hence the direct application of the method of separation of variables will not work.

Now we intend to convert the given problem into two known solvable problems:

one would resemble the problem with homogeneous equation and the other will take care of the nonhomogeneous term.

Seek a solution in the form:

$$u(x, t) = v(x, t) + h(x), \quad (8)$$

where  $h(x)$  is an unknown function of  $x$  alone.

Now, using (8) in (5), we obtain

$$v_{tt} = c^2[v_{xx} + h''(x)] - g. \quad (9)$$

## Finite Vibrating String with Gravity (Contd.)

We select function  $h$  to take care of the non-homogeneous term  $g$  such that

$$c^2 h''(x) = g, \quad (10)$$

and then in turn  $v(x, t)$  satisfies homogeneous wave equation

$$v_{tt} = c^2 v_{xx}. \quad (11)$$

Both functions  $v$  and  $h$  are related by boundary conditions

$$v(0, t) + h(0) = 0, \quad (12a)$$

$$v(L, t) + h(L) = 0, \quad (12b)$$

and initial conditions

$$v(x, 0) + h(x) = \phi(x), \quad (13a)$$

$$v_t(x, 0) = \psi(x). \quad (13b)$$

# Finite Vibrating String with Gravity (Contd.)

Since,  $h$  is a user defined function, we set

$$h(0) = 0 \text{ \& } h(L) = 0. \quad (14)$$

Now the original non-homogeneous problem can be conveniently split into two problems:

## Problem I:

$$\begin{aligned} c^2 h''(x) &= g, \\ h(0) &= 0 = h(L). \end{aligned}$$

## Problem II:

$$\begin{aligned} v_{tt} &= c^2 v_{xx}, \\ v(0, t) &= 0 = v(L, t), \\ v(x, 0) &= \phi(x) - h(x), \quad v_t(x, 0) = \psi(x). \end{aligned}$$

## Finite Vibrating String with Gravity (Contd.)

The solution for Problem I can be easily found by integrating  $h''(x)$  twice:

$$h(x) = \frac{gx^2}{2c^2} + Ax + B.$$

Upon using the conditions  $h(0) = 0 = h(L)$ , we get

$$B = 0 \text{ \& } A = -gL/(2c^2).$$

Hence

$$h(x) = -g \frac{(L-x)x}{2c^2}. \tag{15}$$



## Finite Vibrating String with Gravity (Contd.)

The solution of Problem II is known to us, which is

$$v(x, t) = \sum_{n=1}^{\infty} \sin \frac{n\pi x}{L} \left[ A_n \cos\left(\frac{n\pi ct}{L}\right) + B_n \sin\left(\frac{n\pi ct}{L}\right) \right], \quad (16)$$

where  $A_n$  and  $B_n$  are given, respectively, by

$$A_n = \frac{2}{L} \int_0^L [\phi(x) - h(x)] \sin\left(\frac{n\pi x}{L}\right) dx, \quad n = 1, 2, 3, \dots, \quad (17)$$

$$B_n = \frac{2}{n\pi c} \int_0^L \psi(x) \sin\left(\frac{n\pi x}{L}\right) dx, \quad n = 1, 2, 3, \dots \quad (18)$$

Hence the solution  $u(x, t)$  for our IBVP is given by the sum of (15) and (16).

**Remark:** Clearly, the splitting method would be applicable only when non-homogeneous term is a constant or a function of  $x$ .

# Duhamel's Principle: Finite String Problem

If  $v(x, t, s)$  is the solution of the IBVP

$$v_{tt} - c^2 v_{xx} = 0, \quad (x, t) \in (0, L) \times (0, \infty), \quad (19)$$

$$\text{with BCs } v(0, t, s) = 0, \quad v(L, t, s) = 0, \quad t > 0, \quad s > 0, \quad \text{and} \quad (20)$$

$$\text{with ICs } v(x, 0, s) = 0, \quad v_t(x, 0, s) = f(x, s), \quad s > 0, \quad (21)$$

then  $u(x, t)$  defined by

$$u(x, t) = \int_0^t v(x, t - \tau, \tau) d\tau \quad (22)$$

is the solution to the non-homogeneous problem

$$u_{tt} - c^2 u_{xx} = f(x, t), \quad (x, t) \in (0, L) \times (0, \infty), \quad (23)$$

$$\text{with BCs } u(0, t) = 0, \quad u(L, t) = 0, \quad t > 0, \quad \text{and} \quad (24)$$

$$\text{with ICs } u(x, 0) = 0, \quad u_t(x, 0) = 0. \quad (25)$$

# Finite String Problem: Duhamel's Principle

Example: Find  $u(x, t)$  such that

$$u_{tt} - u_{xx} = t \sin \frac{\pi x}{L}, \quad (x, t) \in (0, L) \times (0, \infty), \quad (26)$$

$$\text{with ICs } u(x, 0) = 0, \quad u_t(x, 0) = 0, \quad x \in (0, L), \quad (27)$$

$$\text{with BCs } u(0, t) = 0, \quad u(L, t) = 0, \quad t > 0. \quad (28)$$

Suppose  $v(x, t, s)$  is a solution to the user defined problem:

$$v_{tt} - v_{xx} = 0, \quad (x, t) \in (0, L) \times (0, \infty), \quad (29)$$

$$\text{with ICs } v(x, 0) = 0, \quad v_t(x, 0) = s \sin \frac{\pi x}{L}, \quad x \in (0, L), \quad s > 0. \quad (30)$$

$$\text{with BCs } v(0, t) = 0, \quad v(L, t) = 0, \quad t > 0. \quad (31)$$

# Finite String Problem: Duhamel's Principle

The solution is given by

$$v(x, t, s) = \sum_{n=1}^{\infty} \sin \frac{n\pi x}{L} \left[ A_n \cos \frac{n\pi t}{L} + B_n \sin \frac{n\pi t}{L} \right] \quad (32)$$

with

$$A_n = \frac{2}{L} \int_0^L \phi(x) \sin \frac{n\pi x}{L} dx = 0, \quad n = 1, 2, 3, \dots$$

$$\begin{aligned} B_n &= \frac{2}{n\pi} \int_0^L \psi(x) \sin \frac{n\pi x}{L} dx, \quad n = 1, 2, 3, \dots \\ &= \frac{2}{n\pi} \int_0^L s \sin \frac{\pi x}{L} \sin \frac{n\pi x}{L} dx. \end{aligned}$$

## Finite String Problem: Duhamel's Principle

Thus,  $B_1 = \frac{sL}{\pi}$  and  $B_n = 0$ ,  $n \neq 1$ , and hence

$$v(x, t, s) = \frac{sL}{\pi} \sin \frac{\pi x}{L} \sin \frac{\pi t}{L}.$$

Then solution  $u(x, t)$  of the given problem is obtained as

$$\begin{aligned} u(x, t) &= \int_0^t v(x, t - \tau, \tau) d\tau \\ &= \int_0^t \frac{\tau L}{\pi} \sin \frac{\pi x}{L} \sin \frac{\pi(t - \tau)}{L} d\tau \\ &= \frac{L}{\pi} \sin \frac{\pi x}{L} \int_0^t \tau \sin \frac{\pi(t - \tau)}{L} d\tau. \end{aligned}$$

# Infinite String Problem: Duhamel's Principle

## Infinite String Problem: Application of Duhamel's Principle

Find  $u(x, t)$  such that

$$u_{tt} - u_{xx} = x - t, \quad (x, t) \in (-\infty, \infty) \times (0, \infty) \quad (33)$$

$$\text{ICs } u(x, 0) = 0, \quad u_t(x, 0) = 0, \quad x \in (-\infty, \infty), \quad s > 0. \quad (34)$$

# Infinite String Problem: Duhamel's Principle

Solution: Suppose  $v(x, t, s)$  solves following user-defined problem

$$v_{tt} - v_{xx} = 0, \quad (x, t) \in (-\infty, \infty) \times (0, \infty) \quad (35)$$

$$\text{ICs } v(x, 0) = 0, \quad v_t(x, 0) = f(x, s) = x - s, \quad x \in (-\infty, \infty), \quad s > 0. \quad (36)$$

D'Alembert's solution is given by

$$v(x, t, s) = \frac{1}{2} \int_{x-t}^{x+t} f(\tau, s) d\tau = \frac{1}{2} \int_{x-t}^{x+t} (\tau - s) d\tau \quad (37)$$

$$= \frac{1}{2} \left[ \frac{\tau^2}{2} - s\tau \right]_{x-t}^{x+t} = xt - ts = t(x - s). \quad (38)$$

Solution to the non-homogeneous problem is given by

$$u(x, t) = \int_0^t v(x, t - \tau, \tau) d\tau \quad (39)$$

$$= \int_0^t 2(t - \tau)(x - \tau) d\tau = -\frac{t^3}{6} + \frac{t^2 x}{2}. \quad (40)$$