



PH101  
Lecture -10  
31st August 2017

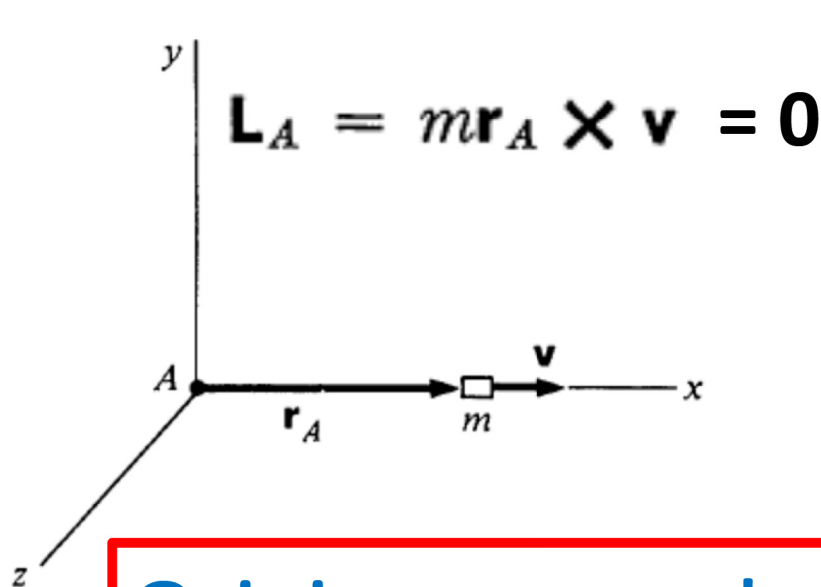
# Angular Momentum & Torque

# Angular Momentum (Moment of Momentum)

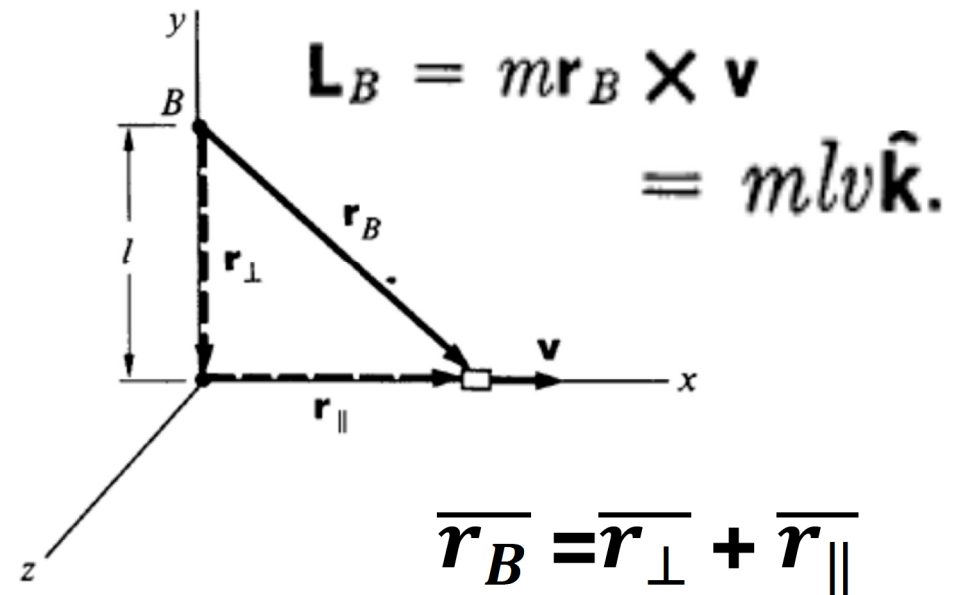
- Angular Momentum of a particle,

$$\mathbf{L} = \mathbf{r} \times \mathbf{p}$$

Direction of  $\mathbf{L}$  is perpendicular to the plane of  $\mathbf{r}$  and  $\mathbf{p}$

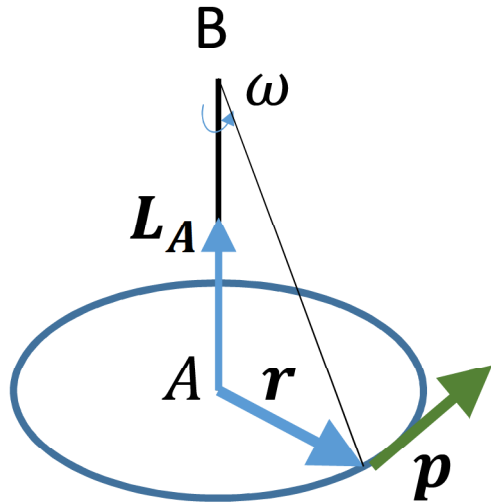


Origin matters!



$$\overline{\mathbf{r}_B} = \overline{\mathbf{r}_\perp} + \overline{\mathbf{r}_\parallel}$$

# Angular Momentum of Conical Pendulum



The pendulum is in steady circular motion with constant angular velocity  $\omega \hat{\mathbf{k}}$ .

Angular momentum about A:

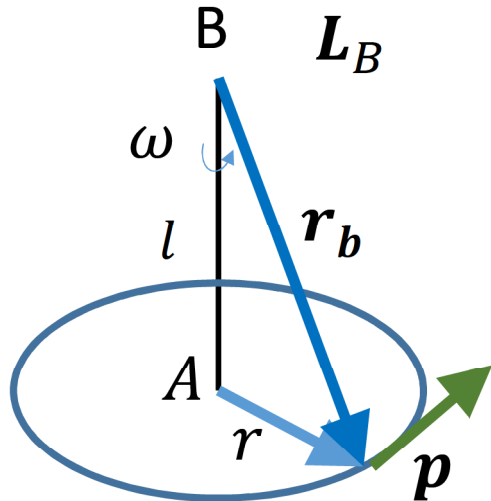
$$\mathbf{L}_A = \mathbf{r} \times \mathbf{p} = rp\hat{\mathbf{k}}$$

$$\mathbf{p} = M\mathbf{v} = Mr\omega$$

$$\mathbf{L}_A = Mr^2\omega \hat{\mathbf{k}}$$

$\mathbf{L}_A$  is **constant** in both magnitude and direction.

# Angular Momentum of Conical Pendulum



Angular momentum about B:

$$\mathbf{L}_B = \mathbf{r}_b \times \mathbf{p}$$

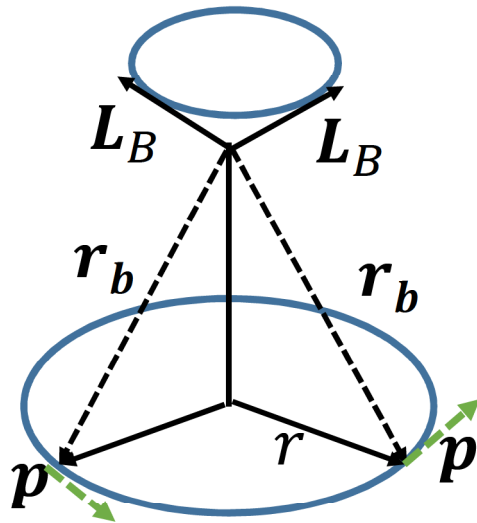
Note that  $\mathbf{r}_b = -l\hat{\mathbf{k}} + r\hat{\mathbf{r}}$  &  $\mathbf{p} = mv\hat{\boldsymbol{\theta}}$

$$\mathbf{L}_B = mv(-l\hat{\mathbf{k}} \times \hat{\boldsymbol{\theta}} + r\hat{\mathbf{r}} \times \hat{\boldsymbol{\theta}})$$

$$= mv(l\hat{\mathbf{r}} + r\hat{\mathbf{k}}) = mr\omega(l\hat{\mathbf{r}} + r\hat{\mathbf{k}})$$

$\mathbf{L}_B$  is constant in magnitude **but direction** is changing because  $\hat{\mathbf{r}}$  is changing!

The z-component of  $\mathbf{L}_B$  ( $mr^2\omega$ ) is constant but the horizontal component ( $mlr\omega$ ) changes its **direction!**



# Torque

For a force,  $F$  acting on a particle at  $r$ , the torque is

$$\boldsymbol{\tau} = \mathbf{r} \times \mathbf{F}$$

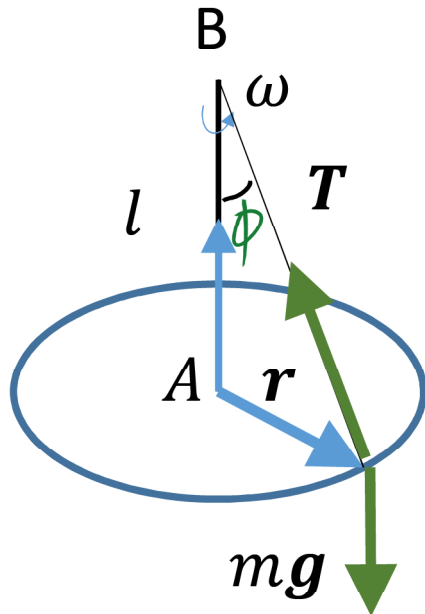
Like angular momentum, the torque depends on the choice of origin.

$$\frac{d\mathbf{L}}{dt} = \frac{d}{dt}(\mathbf{r} \times \mathbf{p})$$

$$= \cancel{\frac{d\mathbf{r}}{dt} \times \mathbf{p}} + \mathbf{r} \times \frac{d\mathbf{p}}{dt}$$

$$\boldsymbol{\tau} = \frac{d\mathbf{L}}{dt} = \mathbf{r} \times \mathbf{F}$$

# Torque on a Conical Pendulum



The pendulum is in steady circular motion with constant angular velocity  $\omega \hat{\mathbf{k}}$ .

Angular momentum about A:

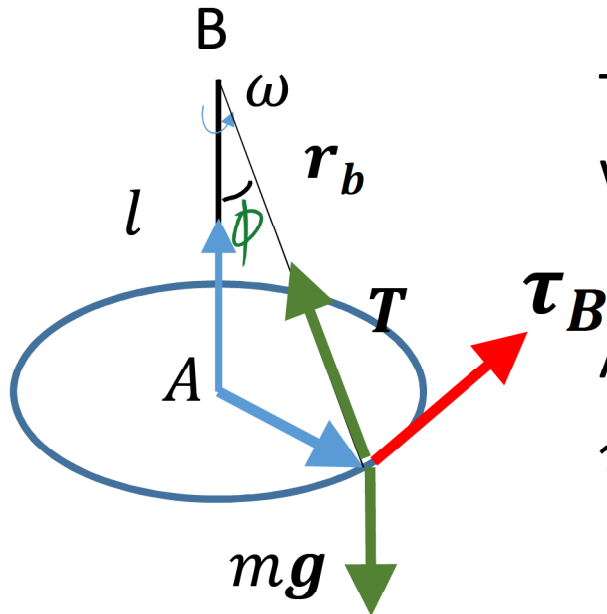
$$\boldsymbol{\tau}_A = \mathbf{r} \times \mathbf{F} = \mathbf{r} \times (\mathbf{T} + mg(-\hat{\mathbf{k}}))$$

$$T \cos\phi - mg = 0 \quad (\text{vertical, } \hat{\mathbf{k}})$$

$$T \sin\phi \quad (\text{horizontal; } -\hat{\mathbf{r}})$$

$$\boldsymbol{\tau}_A = 0 \quad (\mathbf{L}_A \text{ is constant!})$$

# Torque on a Conical Pendulum



The pendulum is in steady circular motion with constant angular velocity  $\omega \hat{\mathbf{k}}$ .

Angular momentum about A:

$$\boldsymbol{\tau}_B = \mathbf{r}_b \times \mathbf{F} = \mathbf{r}_b \times \mathbf{T} + \mathbf{r}_b \times mg(-\hat{\mathbf{k}})$$

$$\mathbf{r}_b = -l\hat{\mathbf{k}} + r\hat{\mathbf{r}}$$

$$\boldsymbol{\tau}_B = -l\hat{\mathbf{k}} \times mg(-\hat{\mathbf{k}}) + r\hat{\mathbf{r}} \times mg(-\hat{\mathbf{k}})$$

$$\hat{\mathbf{r}}, \hat{\boldsymbol{\theta}}, \hat{\mathbf{k}}$$

$$\boldsymbol{\tau}_B = mgr \hat{\boldsymbol{\theta}} = \frac{dL_B}{dt} \quad (\boldsymbol{\tau}_B \text{ changes its direction!})$$

# Torque of Conical Pendulum

Angular momentum about B:

$$\mathbf{L}_B = mr\omega(l\hat{\mathbf{r}} + r\hat{\mathbf{k}})$$

$$\frac{d\mathbf{L}_B}{dt} = mr\omega l \frac{d\hat{\mathbf{r}}}{dt} = mrl\omega\omega\hat{\boldsymbol{\theta}}$$

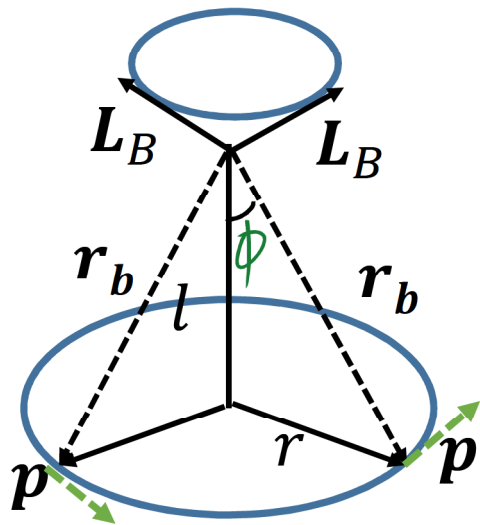
$$T \cos\phi = mg \quad (\text{vertical, } \hat{\mathbf{k}})$$

$$T \sin\phi = mr\omega^2 \quad (\text{horizontal; } -\hat{\mathbf{r}})$$

$$g/r\omega^2 = \cot\phi = l/r$$

$$\omega^2 = g/l$$

$$\frac{d\mathbf{L}_B}{dt} = mrl\omega^2\hat{\boldsymbol{\theta}} = mgr\hat{\boldsymbol{\theta}} = \boldsymbol{\tau}_B$$



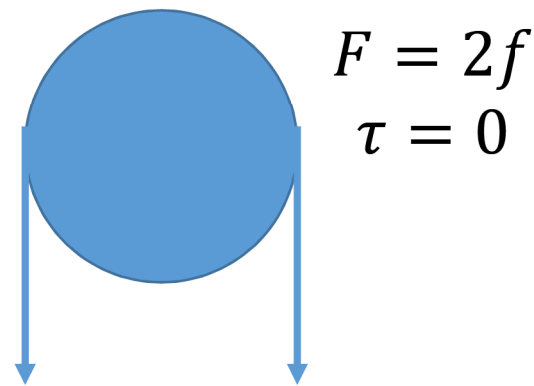
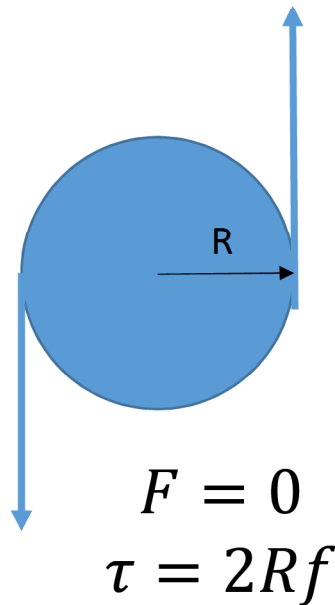


# Force & Torque

Torque and force are very different quantities (KK)

There can be a torque on a system with zero net force

And there can be net force on a system with zero net torque.



# Rigid Bodies

- Many particle system
- Interparticle distance for all pairs of particles  $|r_{ij}| = c_{ij}$
- Retains shape and size
- 6 degrees of freedom

## FIXED AXIS ROTATION

- All particles rotate about an axis fixed in space
- All particles have the same angular speed  $\omega$
- Let us choose the z-axis as the axis of rotation, with  $\boldsymbol{\omega} = \omega \hat{k}$

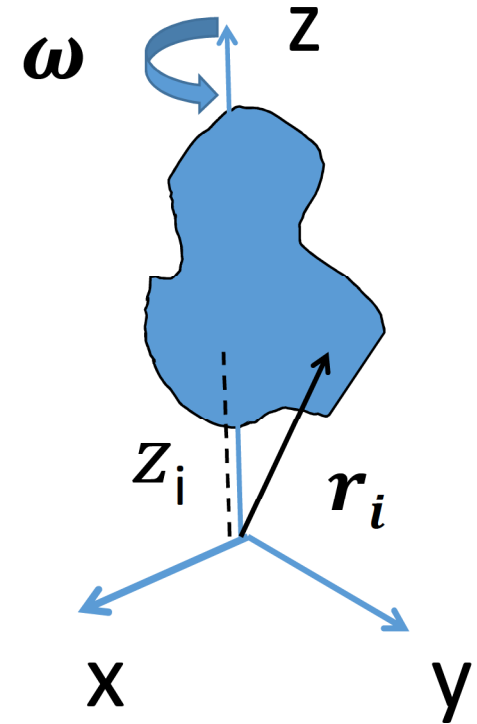
# Rotation about a Fixed Axis

$$\mathbf{L} = \sum_i \mathbf{r}_i \times \mathbf{p}_i$$

$$\mathbf{L} = \sum_i m_i \mathbf{r}_i \times (\boldsymbol{\omega} \times \mathbf{r}_i) = \sum_i m_i r_i^2 \boldsymbol{\omega} - \sum_i m_i (\mathbf{r}_i \cdot \boldsymbol{\omega}) \mathbf{r}_i$$

$$L_z = \sum_i m_i (r_i^2 - z_i^2) \omega \hat{\mathbf{k}} = \sum_i m_i (x_i^2 + y_i^2) \omega \hat{\mathbf{k}}$$

$$L_z = I_{zz} \omega \hat{\mathbf{k}}$$



# Moment of Inertia

- $I_{zz}$  is called the moment of inertia about the z-axis
- $\sum_i m_i (x_i^2 + y_i^2) = I_{zz}$  is a purely geometric quantity. Does not depend on the motion.
- For continuous bodies

$$I_{zz} = \int dm(x^2 + y^2)$$

$$I_{zz} = \int \rho(x^2 + y^2)dV$$

Where  $\rho(\mathbf{r})$  is the density and  $dV$  is the volume element.

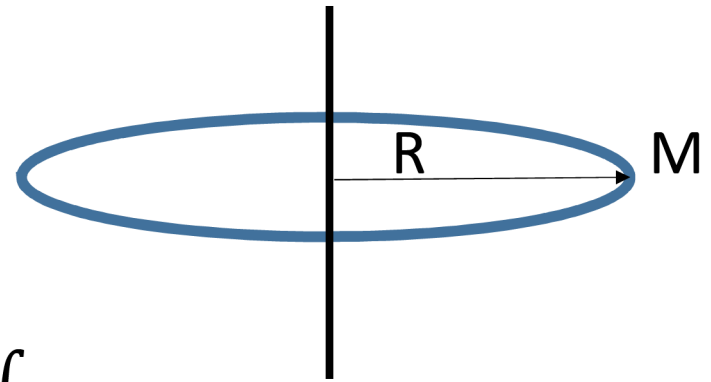
## Example 1

A hoop/ring of mass  $M$  and radius  $R$ .

$$I_{ZZ} = \int dm(x^2 + y^2) = \int dm R^2 = MR^2$$

$$\text{Or, } I_{ZZ} = \int (\lambda R d\theta) R^2 = R^3 2\pi \lambda = MR^2$$

$$\lambda = M/2\pi R$$

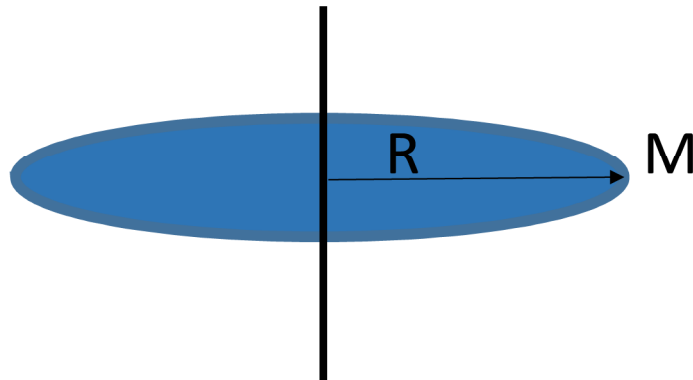


## Example 2

A thin disk of mass  $M$  and radius  $R$ .

$$I_{ZZ} = \int \sigma \, dx \, dy \, (x^2 + y^2) = \int r^2 \, \sigma \, dr \, r \, d\theta = \frac{R^4}{4} 2\pi \, \sigma = \frac{MR^2}{2}$$

$$\sigma = \frac{M}{\pi R^2}$$

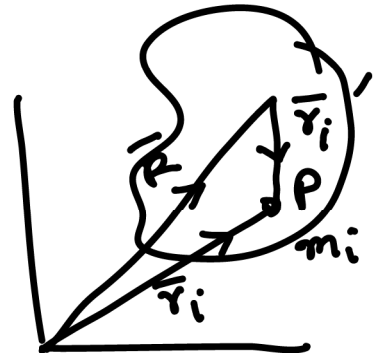


# System of Particles (General Motion)

$$\vec{L} = \sum_{i=1} \vec{r}_i \times m_i \dot{\vec{r}}_i$$

$$\vec{R} = \frac{\sum m_i \vec{r}_i}{M} \quad (CM)$$

$$\left| \begin{array}{l} \vec{r}_i = \vec{R} + \vec{r}_i' \\ \dot{\vec{r}}_i = \dot{\vec{R}} + \dot{\vec{r}}_i' \end{array} \right.$$



$$\vec{L} = \sum m_i (\vec{R} + \vec{r}_i') \times (\dot{\vec{R}} + \dot{\vec{r}}_i')$$

$$= \sum m_i (\vec{R} \times \dot{\vec{R}}) + \sum m_i \vec{R} \times \dot{\vec{r}}_i' + \sum m_i \vec{r}_i' \times \dot{\vec{R}} + \sum m_i \vec{r}_i' \times \dot{\vec{r}}_i'$$

$$= \underbrace{\vec{R} \times M \dot{\vec{R}}}_{\vec{L}_0} + \vec{R} \times \sum m_i \dot{\vec{r}}_i' + (\sum m_i \vec{r}_i') \times \dot{\vec{R}} + \underbrace{\sum \vec{r}_i' \times m_i \dot{\vec{r}}_i'}_{\vec{L}_{CM}}$$

But

$$M\bar{R} = \sum m_i \bar{r}_i = \sum m_i (\bar{R} + \bar{r}_i') = M\bar{R} + \sum m_i \bar{r}_i' \quad \text{ie } \underline{\underline{\sum m_i \bar{r}_i' = 0}}$$

So

$$\bar{L} = \bar{R} \times M\bar{V}_{cm} + \sum_i \bar{r}_i' \times \bar{p}_i'$$

$$\text{Or } \underline{\underline{\bar{L} = \bar{L}_0 + \bar{L}_{cm}}}$$

$$\bar{p}_i' = m_i \dot{\bar{r}}_i' = m_i \bar{v}_i' \\ \text{w.r. to CM}$$

Note:  $\bar{L}_0$  is the  $\bar{L}$  of the Center of Mass.

$$\text{ie, } \bar{L}_0 = \bar{R} \times M\bar{V}_{cm} = \bar{R} \times \bar{p}_{cm}$$

$\bar{L}_{cm}$  is the  $\underline{\underline{\sum \bar{r}_i' \times \bar{p}_i'}}$  with respect to the CM!



Torque

$$\bar{\tau} = \sum \bar{r}_i \times \bar{f}_i$$

$$\left| \bar{f}_i = \bar{f}_i^{\text{ext}} + \bar{f}_i^{\text{int}} \text{ on } i^{\text{th}} \text{ particle} \right.$$

$$= \sum (\bar{R} + \bar{r}_i') \times \bar{f}_i$$

$$= \bar{R} \times \sum_i \bar{f}_i + \sum \bar{r}_i' \times \bar{f}_i$$

$$= \bar{R} \times \left( \sum_i \bar{f}_i^{\text{int}} + \sum_i \bar{f}_i^{\text{ext}} \right) + \sum \bar{r}_i' \times \bar{f}_i$$

But  $\sum_i \bar{f}_i^{\text{int}} = 0!$

$$\bar{\tau} = \bar{R} \times \bar{F} + \sum \bar{r}_i' \times \bar{f}_i^{\text{int}} + \sum \bar{r}_i' \times \bar{f}_i^{\text{ext}}$$

$\bar{F}$  - is the total external force on the body!

$$\sum \bar{r}_i' \times \bar{f}_i^{\text{int}} \quad \Bigg| \quad \bar{f}_i^{\text{int}} = \sum_{j \neq i} \bar{f}_{ij} \quad \left( \begin{array}{l} \text{force on } i \text{ due} \\ \text{to } j \end{array} \right)$$

has terms like  $\bar{r}_i' \times \bar{f}_{ij} + \bar{r}_j' \times \bar{f}_{ji} = (\bar{r}_i' - \bar{r}_j') \times \bar{f}_{ij}$

$$\bar{f}_{ji} = -\bar{f}_{ij}$$

Now if the internal forces between particles  $i$  &  $j$  act along the line connecting them  $(\bar{r}_i' - \bar{r}_j')$  then  $(\bar{r}_i' - \bar{r}_j') \times \bar{f}_{ij} = 0$

So contribution to  $\bar{\tau}$  due to internal forces vanish!

So 
$$\underline{\underline{\bar{\tau} = \bar{R} \times \bar{F} + \sum \bar{r}_i' \times \bar{f}_i^{\text{ext}}}}$$