

1. Prove the mean value theorem: Let  $P$  be an interior point of a volume  $V$ . Let  $y$  be a solution of the Laplace equation in  $V$ . Then  $y(P)$  is the average of  $y$  over the surface of any sphere in  $V$  centered about  $P$ . [Hint: Use the integral equation.] Prove that the solution of the Laplace equation cannot have a maximum or a minimum in  $V$ .
2. Consider the Laplace equation  $\nabla^2\phi = 0$  in a volume  $V$  with boundary  $S$ .

(a) Prove using the Green's identity, that for a function  $f$ ,

$$\int_V (f\nabla^2 f + |\nabla f|^2) dv = \oint_S f (\nabla f \cdot \hat{\mathbf{n}}) dS.$$

(b) Prove that the solution (assuming that it exists) to the Laplace equation in  $V$  with either Dirichlet or Neumann boundary conditions must be unique.

3. Prove that the Green's function for Laplace equation must be symmetric under exchange of its arguments, that is,  $G(\mathbf{r}, \mathbf{r}') = G(\mathbf{r}', \mathbf{r})$ . [Note: This result is true for all self-adjoint operators. Try and prove this result for Green's functions of Sturm-Liouville equations.]
4. Show that the Dirichlet Green's function for the unbounded space between  $z = 0$  and  $z = L$  planes is given by

$$G(\mathbf{r}, \mathbf{r}') = \frac{4}{L} \sum_{n=1}^{\infty} \sum_{m=-\infty}^{\infty} e^{im(\phi-\phi')} \sin\left(\frac{n\pi z}{L}\right) \sin\left(\frac{n\pi z'}{L}\right) I_m\left(\frac{n\pi}{L}\rho_{<}\right) K_m\left(\frac{n\pi}{L}\rho_{>}\right).$$

5. For the same geometry of problem 4, show that alternate form of the Green's function is

$$G(\mathbf{r}, \mathbf{r}') = 2 \sum_{m=-\infty}^{\infty} \int_0^k dk e^{im(\phi-\phi')} J_m(k\rho) J_m(k\rho') \frac{\sinh(kz_{<}) \sinh[k(L-z_{>})]}{\sinh(kL)}.$$

6. Consider two parallel conducting plates  $z = 0$  and  $z = L$ . The potential on the  $z = 0$  plate is zero, and on  $z = L$  is given by

$$\Phi(\rho, \phi, L) = \begin{cases} V & \rho \leq a \\ 0 & \rho > a. \end{cases}$$

(a) Show that the potential between the plates can be written as

$$\Phi(\rho, \phi, z) = V \int_0^{\infty} d\lambda J_1(\lambda) J_0(\lambda\rho/a) \frac{\sinh(\lambda z/a)}{\sinh(\lambda L/a)}$$

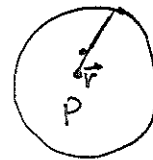
## Tutorial 7

Q1. (a) Given, for a volume  $V$ ,

$$\phi(\vec{r}) = \frac{1}{4\pi} \oint_S \left[ \frac{1}{R} \frac{\partial \phi}{\partial n'} - \phi \frac{\partial}{\partial n'} \left( \frac{1}{R} \right) \right] ds'$$

if  $\phi$  is a sol<sup>n</sup> of the Laplace Eq.

let  $S$  be a spherical shell of radius  $a$  about  $\vec{r}$ . Then  $\frac{1}{R} = \frac{1}{|\vec{r} - \vec{r}'|}$



$$\therefore \frac{1}{R} = \frac{1}{a} \text{ on } S \text{ and } \hat{n} = \widehat{(\vec{r} - \vec{r}')}$$

$$\hat{n} \cdot \nabla' \left( \frac{1}{|\vec{r} - \vec{r}'|} \right) = - \widehat{(\vec{r} - \vec{r}')} \cdot \frac{(\vec{r} - \vec{r}')}{a^3} = -\frac{1}{a^2}$$

and 
$$\oint_S \nabla' \phi \cdot \hat{n}' ds = \int_V \nabla'^2 \phi ds = 0.$$

$$\therefore \phi(\vec{r}) = \frac{1}{4\pi a^2} \oint_S \phi ds'.$$

(b) If  $\phi$  has a maximum at  $P$  then for some  $\epsilon$ -nbd of  $P$

$$\phi(\vec{r}) \leq \phi(P). \text{ Thus on } \epsilon\text{-ball}$$

$$\frac{1}{4\pi a^2} \int \phi ds < \phi(P)$$

↑  
Strictly less than

This is in contrast to previous Result.

Q2. (a) Now by divergence theorem  $\int_V \nabla \cdot A dv = \oint_S A \cdot \hat{n} ds$

Choose  $A = f \nabla f$ ,  $\nabla \cdot A = \nabla \cdot (f \nabla f) = f \nabla^2 f + \nabla f \cdot \nabla f$   
 $= f \nabla^2 f + |\nabla f|^2$

$$\therefore \int_V (f \nabla^2 f + |\nabla f|^2) dv = \oint_S f \frac{\partial f}{\partial n} ds \quad \text{--- (1)}$$

(b) if  $\phi_1$  is sol<sup>n</sup> of  $\nabla^2 \phi_1 = 0 \Rightarrow$  with DBC  $\phi_1(P) = \eta(P)$  on  $S$

So if  $\phi_2$ , i.e.  $\nabla^2 \phi_2 = 0$  . "  $\phi_2(P) = \eta(P)$  on  $S$

This implies that  $\dot{\phi} = \phi_1 - \phi_2$  satisfies

$$\nabla^2 \phi = 0 \quad \text{s.t.} \quad \phi = 0 \quad \text{on } S.$$

Thus if  $f = \phi$  in ①

$$\int_V (\nabla \phi)^2 dv = 0.$$

$$\Rightarrow \nabla \phi = 0 \quad \text{on } V$$

$$\Rightarrow \phi = \text{const on } V. \quad \text{but the const must be zero since } \phi = 0 \quad \text{on } S.$$

$$\Rightarrow \phi_1 = \phi_2 \quad \text{not } \neq.$$

Q3. let  $r_1, r_2$  be two points in  $V$  bounded by  $S$ . Consider Dirichlet Green's function.

$$\nabla^2 G(\bar{r}, \bar{r}') = -4\pi \delta(\bar{r} - \bar{r}') \quad \text{and } G(\bar{r}, \bar{r}') = 0 \quad \text{on } S.$$

Then

$$\begin{aligned} G(r_2, r_1) \nabla^2 G(r, r_2) - G(r, r_2) \nabla^2 G(r, r_1) \\ = [G(r, r_1) \delta(r - r_2) - G(r, r_2) \delta(r - r_1)] (-4\pi) \end{aligned}$$

Integrate on Surface.

$$\begin{aligned} \text{RHS} &= \int_V [G \nabla^2 G(r, r_2) - G(r, r_2) \nabla^2 G(r, r_1)] dv \\ &= \int_V \nabla \cdot [G(r, r_1) \nabla G(r, r_2) - G(r, r_2) \nabla G(r, r_1)] dv \\ &= \oint_S [G(r, r_1) \nabla G(r, r_2) - G(r, r_2) \nabla G(r, r_1)] ds \end{aligned}$$

$$= 0 \quad \text{since } G = 0 \quad \text{on } S \quad [\text{clearly will apply to Neumann case too.}]$$

$$\text{RHS} = G(r_2, r_1) - G(r_1, r_2)$$

$$\Rightarrow G(r_1, r_2) = G(r_2, r_1).$$

Q4: Assume that

$$G(\vec{r}, \vec{r}') = \frac{4}{L} \sum_{m,n} g_m(r, r') \sin\left(\frac{n\pi z}{L}\right) \sin\left(\frac{n\pi z'}{L}\right) e^{-im(\phi - \phi')}$$

Then  $G = 0$  if  $z = 0$  or  $z = L$ .

$$\text{Now, } \nabla^2 G = \frac{4}{L} \sum_{m,n} \left[ \frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial g}{\partial r} \right) - \left( \frac{m^2}{r^2} + k^2 \right) g \right] e^{im(\phi - \phi')} \\ \times \sin\left(\frac{n\pi z}{L}\right) \sin\left(\frac{n\pi z'}{L}\right)$$

$$= -4\pi \delta^3(\vec{r} - \vec{r}') = -\frac{4\pi}{r} \delta(r - r') \delta(\phi - \phi') \delta(z - z')$$

$$\text{Thus if } \left[ \frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial}{\partial r} \right) - \left( \frac{m^2}{r^2} + k^2 \right) \right] g = -\frac{1}{r} \delta(r - r')$$

$$\text{then } \nabla^2 G = -\frac{4\pi}{\pi L} \delta(r - r') \sum_m e^{im(\phi - \phi')} \sum_n \sin\left(\frac{n\pi z}{L}\right) \sin\left(\frac{n\pi z'}{L}\right) \\ = -\frac{4\pi}{\pi L r} \delta(r - r') \cdot 2\pi \delta(\phi - \phi') \cdot \frac{L}{2} \delta(z - z') \\ = -\frac{4\pi}{r} \delta(r - r') \delta(\phi - \phi') \delta(z - z') = -4\pi \delta^3(\vec{r} - \vec{r}')$$

The Eq. for  $g_m$  is then  $\rho = kr$

$$\frac{1}{\rho} \frac{\partial}{\partial \rho} \left( \rho \frac{\partial g}{\partial \rho} \right) - \left( \frac{m^2}{\rho^2} + k^2 \right) g = -\delta(\rho - \rho')$$

$$\therefore g_1 = I_m(kr) \quad , \quad g_2 = K_m(kr) \\ g_1(0) = 0 \quad , \quad g_2(\infty) = 0$$

From properties of Bessel fns.

$$W(g_1, g_2) = -\frac{1}{kr} = -\frac{1}{\rho}$$

$$C = -\frac{1}{\rho(e') W(e')} = 1$$

$$\therefore g_{m,n}(\rho) = I_m(kr_c) K_m(kr_r)$$

$$\therefore G(\vec{r}, \vec{r}') = \frac{4}{L} \sum_{m,n} \frac{1}{\rho} e^{im(\phi - \phi')} \cdot \sin\left(\frac{n\pi z}{L}\right) \sin\left(\frac{n\pi z'}{L}\right) \cdot I_m\left(\frac{n\pi}{L} r_c\right) K_m\left(\frac{n\pi}{L} r_r\right)$$

Q5. Orthogonality of Bessel Fns.

$$\int_0^{\infty} J_m(kr) J_m(k'r) r dr = \frac{1}{k} \delta(k-k')$$

$$\therefore \delta(r-r') = \int_0^{\infty} A_m(k) J_m(kr) dk$$

$$\therefore \int r J_m(k'r) \delta(r-r') dk = \int_0^{\infty} A_m(k) \int_0^{\infty} r J_m(k'r) J_m(kr) dr dk$$

$$r' J_m(k'r') = \frac{1}{k'} A_m(k')$$

$$\therefore A_m(k') = \frac{1}{k'} (k'r') J_m(k'r')$$

$$\Rightarrow \frac{\delta(r-r')}{r'} = \int_0^{\infty} k J_m(kr) J_m(kr') dk$$

$$\therefore G(r, r') = 2 \sum_{m=-\infty}^{\infty} \int_0^{\infty} dk J_m(kr) J_m(kr') \cdot e^{im(\phi-\phi')} \cdot g_{m,k}(z, z')$$

$$\therefore \nabla^2 G = 2 \sum_m \int dk \left[ \frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2}{\partial \theta^2} + \frac{\partial^2}{\partial z^2} \right] J_m(kr) J_m(kr') e^{im(\phi-\phi')} g_{m,k}(z, z')$$

$$= 2 \sum_m \int dk \left[ \left( \frac{m^2}{r^2} - k^2 \right) - \frac{1}{r^2} m^2 + \frac{\partial^2}{\partial z^2} \right] J_m(kr) J_m(kr') e^{im(\phi-\phi')} g_{m,k}(z-z')$$

$$\therefore \text{if } \left( \frac{\partial^2}{\partial z^2} - k^2 \right) g_{m,k}(z, z') = -\frac{k}{z} \delta(z-z')$$

$$\begin{aligned} \text{then } \nabla^2 G &= -2 \sum_m \int dk \delta(z-z') J_m(kr) J_m(kr') e^{im(\phi-\phi')} \\ &= -2 \cdot \delta(z-z') \cdot 2\pi \delta(\phi-\phi') \cdot \frac{\delta(r-r')}{r'} = -4\pi \delta^3(\vec{r}-\vec{r}') \end{aligned}$$

$$g_m(z_1, z_2) = \frac{\sinh(kz_1) \sinh(k(L-z_2))}{\sinh(kL)}$$

Q5. Do the same as in Q4 except put  $\frac{d^2}{dz^2} Q(z) = +k^2 Q(z)$ .

Q6. Using Result of Q5.

$$\therefore \phi(r, \phi, z) = \int \frac{\partial G(r, r')}{\partial n} \phi(r, \phi, \theta) r dr d\theta$$

Since  $\int_0^{2\pi} e^{im(\theta-\theta')} d\theta = \begin{cases} 2\pi & \text{if } m=0 \\ 0 & \text{otherwise} \end{cases}$

$$\therefore \phi(r, \phi, z) = \int_0^\infty dk V \int_0^a z dr J_0(kr) J_0(kr')$$

$$\hat{n} = +k \Rightarrow \frac{\partial G}{\partial n} = \left( \right) \frac{\cosh(kz) \cosh(kL-z')}{\sinh(kL)}$$

$$= - \left( \right) \frac{\sinh(kz')}{\sinh(kL)}$$

$$\Rightarrow \phi(r, \phi, z) = \int_0^\infty dk V \frac{\sinh(kz')}{\sinh(kL)} J_0(kr') \int z dz J_0(kr)$$

$$= \int_0^\infty \frac{2V \sinh(kz')}{\sinh(kL)} J_0(kr') J_0(kr) dz$$

$\Rightarrow$  Result.