

[Note: Only the first five problems will be discussed in tutorial classes. You should attempt the remaining problems and in case you get stuck, ask your tutorial instructor]

1. [G 2.34] Consider two concentric spherical shells, of radii a and b . Suppose the inner one carries a charge q , and the outer one a charge $-q$ (both of them uniformly distributed over the surface). Calculate the energy of this configuration, (a) using $W = \frac{\epsilon_0}{2} \int E^2 d\tau$, and (b) using $W_{tot} = W_1 + W_2 + \epsilon_0 \int \mathbf{E}_1 \cdot \mathbf{E}_2 d\tau$.

(a) $W = \frac{\epsilon_0}{2} \int E^2 d\tau$. $\mathbf{E} = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{\mathbf{r}}$ ($a < r < b$), zero elsewhere.

$$W = \frac{\epsilon_0}{2} \left(\frac{q}{4\pi\epsilon_0} \right)^2 \int_a^b \left(\frac{1}{r^2} \right)^2 4\pi r^2 dr = \frac{q^2}{8\pi\epsilon_0} \int_a^b \frac{1}{r^2} dr = \frac{q^2}{8\pi\epsilon_0} \left(\frac{1}{a} - \frac{1}{b} \right).$$

(b) $W_1 = \frac{1}{8\pi\epsilon_0} \frac{q^2}{a}$, $W_2 = \frac{1}{8\pi\epsilon_0} \frac{q^2}{b}$, $\mathbf{E}_1 = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{\mathbf{r}}$ ($r > a$), $\mathbf{E}_2 = \frac{1}{4\pi\epsilon_0} \frac{-q}{r^2} \hat{\mathbf{r}}$ ($r > b$). So

$$\mathbf{E}_1 \cdot \mathbf{E}_2 = \left(\frac{1}{4\pi\epsilon_0} \right)^2 \frac{-q^2}{r^4}, \quad (r > b), \text{ and hence } \int \mathbf{E}_1 \cdot \mathbf{E}_2 d\tau = - \left(\frac{1}{4\pi\epsilon_0} \right)^2 q^2 \int_b^\infty \frac{1}{r^4} 4\pi r^2 dr = - \frac{q^2}{4\pi\epsilon_0 b}.$$

$$W_{tot} = W_1 + W_2 + \epsilon_0 \int \mathbf{E}_1 \cdot \mathbf{E}_2 d\tau = \frac{1}{8\pi\epsilon_0} q^2 \left(\frac{1}{a} + \frac{1}{b} - \frac{2}{b} \right) = \frac{q^2}{8\pi\epsilon_0} \left(\frac{1}{a} - \frac{1}{b} \right).$$

2. [G 2.45] A sphere of radius R carries a charge density $\rho(r) = kr$ (where k is a constant). Find the energy of the configuration. Check your answer by calculating it in at least two different ways.

First let's determine the electric field inside and outside the sphere, using Gauss's law:

$$\epsilon_0 \oint \mathbf{E} \cdot d\mathbf{a} = \epsilon_0 4\pi r^2 E = Q_{enc} = \int \rho d\tau = \int (k\bar{r}) \bar{r}^2 \sin\theta d\bar{r} d\theta d\phi = 4\pi k \int_0^r \bar{r}^3 d\bar{r} = \begin{cases} \pi k r^4 & (r < R), \\ \pi k R^4 & (r > R). \end{cases}$$

$$\text{So } \mathbf{E} = \frac{k}{4\epsilon_0} r^2 \hat{\mathbf{r}} \quad (r < R); \quad \mathbf{E} = \frac{kR^4}{4\epsilon_0 r^2} \hat{\mathbf{r}} \quad (r > R).$$

Method I:

$$\begin{aligned} W &= \frac{\epsilon_0}{2} \int E^2 d\tau \quad (\text{Eq. 2.45}) = \frac{\epsilon_0}{2} \int_0^R \left(\frac{kr^2}{4\epsilon_0} \right)^2 4\pi r^2 dr + \frac{\epsilon_0}{2} \int_R^\infty \left(\frac{kR^4}{4\epsilon_0 r^2} \right)^2 4\pi r^2 dr \\ &= 4\pi \frac{\epsilon_0}{2} \left(\frac{k}{4\epsilon_0} \right)^2 \left\{ \int_0^R r^6 dr + R^8 \int_R^\infty \frac{1}{r^2} dr \right\} = \frac{\pi k^2}{8\epsilon_0} \left\{ \frac{R^7}{7} + R^8 \left(-\frac{1}{r} \right) \Big|_R^\infty \right\} = \frac{\pi k^2}{8\epsilon_0} \left(\frac{R^7}{7} + R^7 \right) \\ &= \boxed{\frac{\pi k^2 R^7}{7\epsilon_0}}. \end{aligned}$$

Method II:

$$\begin{aligned} W &= \frac{1}{2} \int \rho V d\tau \quad (\text{Eq. 2.43}). \\ \text{For } r < R, \quad V(r) &= - \int_\infty^r \mathbf{E} \cdot d\mathbf{l} = - \int_\infty^R \left(\frac{kR^4}{4\epsilon_0 r^2} \right) dr - \int_R^r \left(\frac{kr^2}{4\epsilon_0} \right) dr = - \frac{k}{4\epsilon_0} \left\{ R^4 \left(-\frac{1}{r} \right) \Big|_\infty^R + \frac{r^3}{3} \Big|_R^r \right\} \\ &= - \frac{k}{4\epsilon_0} \left(-R^3 + \frac{r^3}{3} - \frac{R^3}{3} \right) = \frac{k}{3\epsilon_0} \left(R^3 - \frac{r^3}{4} \right). \\ \therefore W &= \frac{1}{2} \int_0^R (kr) \left[\frac{k}{3\epsilon_0} \left(R^3 - \frac{r^3}{4} \right) \right] 4\pi r^2 dr = \frac{2\pi k^2}{3\epsilon_0} \int_0^R \left(R^3 r^3 - \frac{1}{4} r^6 \right) dr \\ &= \frac{2\pi k^2}{3\epsilon_0} \left\{ R^3 \frac{R^4}{4} - \frac{1}{4} \frac{R^7}{7} \right\} = \frac{\pi k^2 R^7}{2 \cdot 3\epsilon_0} \left(\frac{6}{7} \right) = \frac{\pi k^2 R^7}{7\epsilon_0}. \end{aligned}$$

3. [G 2.36] Two spherical cavities, of radii a and b , are hollowed out from the interior of a (neutral) conducting sphere of radius R (Fig.). At the center of each cavity a point charge is placed - call these charges q_a and q_b .

- Find the surface charges σ_a , σ_b , and σ_R .
- What is the field outside the conductor?
- What is the field within each cavity?
- What is the force on q_a and q_b ?
- Which of these answers would change if a third charge, q_c , were brought near the conductor?

(a) $\sigma_a = -\frac{q_a}{4\pi a^2}$; $\sigma_b = -\frac{q_b}{4\pi b^2}$; $\sigma_R = \frac{q_a + q_b}{4\pi R^2}$.

(b) $\mathbf{E}_{\text{out}} = \frac{1}{4\pi\epsilon_0} \frac{q_a + q_b}{r^2} \hat{\mathbf{r}}$, where \mathbf{r} = vector from center of large sphere.

(c) $\mathbf{E}_a = \frac{1}{4\pi\epsilon_0} \frac{q_a}{r_a^2} \hat{\mathbf{r}}_a$, $\mathbf{E}_b = \frac{1}{4\pi\epsilon_0} \frac{q_b}{r_b^2} \hat{\mathbf{r}}_b$, where \mathbf{r}_a (\mathbf{r}_b) is the vector from center of cavity a (b).

(d) Zero.

(e) σ_R changes (but not σ_a or σ_b); $\mathbf{E}_{\text{outside}}$ changes (but not \mathbf{E}_a or \mathbf{E}_b); force on q_a and q_b still zero.

4. [G 2.43] Find the net force that the southern hemisphere of a uniformly charged sphere exerts on the northern hemisphere. Express your answer in terms of the radius R and the total charge Q .

From Prob. 2.12, the field inside a uniformly charged sphere is: $\mathbf{E} = \frac{1}{4\pi\epsilon_0} \frac{Q}{R^3} \mathbf{r}$. So the force per unit volume is $\mathbf{f} = \rho \mathbf{E} = \left(\frac{Q}{\frac{4}{3}\pi R^3}\right) \left(\frac{Q}{4\pi\epsilon_0 R^3}\right) \mathbf{r} = \frac{3}{\epsilon_0} \left(\frac{Q}{4\pi R^3}\right)^2 \mathbf{r}$, and the force in the z direction on $d\tau$ is:

$$dF_z = f_z d\tau = \frac{3}{\epsilon_0} \left(\frac{Q}{4\pi R^3}\right)^2 r \cos\theta (r^2 \sin\theta dr d\theta d\phi).$$

The total force on the "northern" hemisphere is:

$$F_z = \int f_z d\tau = \frac{3}{\epsilon_0} \left(\frac{Q}{4\pi R^3}\right)^2 \int_0^R r^3 dr \int_0^{\pi/2} \cos\theta \sin\theta d\theta \int_0^{2\pi} d\phi$$

$$= \frac{3}{\epsilon_0} \left(\frac{Q}{4\pi R^3}\right)^2 \left(\frac{R^4}{4}\right) \left(\frac{\sin^2\theta}{2} \Big|_0^{\pi/2}\right) (2\pi) = \boxed{\frac{3Q^2}{64\pi\epsilon_0 R^2}}.$$

5. Let $q_1 (> 0)$ and $-q_2$, (where $q_2 > 0$) be two charges located at $(d, 0, 0)$ and $(-d, 0, 0)$ respectively. Show that the zero potential surface is spherical with center on x axis. If center of the surface is at $(c, 0, 0)$ and radius is R , find c/d and R/d in terms of q_1/q_2 . Sketch surfaces for $q_1/q_2 = \frac{1}{4}, \frac{1}{3}, \frac{1}{2}, 1, 2, 3, 4$.

If a point (x, y, z) is on $V = 0$ surface, then

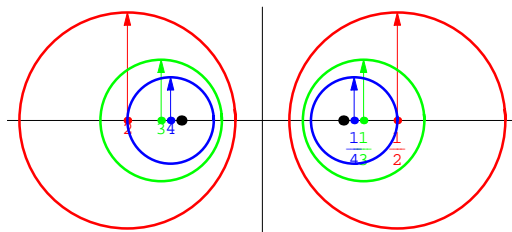
$$\frac{q_1}{((x-d)^2 + y^2 + z^2)} = \frac{q_2}{((x+d)^2 + y^2 + z^2)}$$

$$\therefore \left[x - d \frac{1+\alpha^2}{1-\alpha^2} \right]^2 + y^2 + z^2 = d^2 \frac{4\alpha^2}{|1-\alpha^2|^2}$$

where, $\alpha = q_1/q_2$. This is an equation of a sphere with centre

$$c = d \frac{1+\alpha^2}{1-\alpha^2} \quad \text{and radius} \quad R = d \frac{2\alpha}{|1-\alpha^2|}$$

Figure shows surfaces for $\alpha = \frac{1}{4}, \frac{1}{3}, \frac{1}{2}, 1, 2, 3, 4$. As the ratio $\alpha \rightarrow 1$, $R \rightarrow \infty$ and $c \rightarrow \pm\infty$ and equipotential surface approaches yz plane.



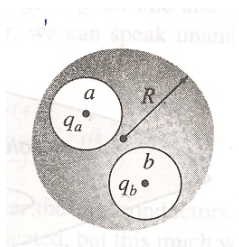
6. [G 2.35] A metal sphere of radius R , carrying charge q , is surrounded by a thick concentric metal shell (inner radius a , outer radius b , as shown in Fig.). The shell carries no net charge.

- Find the surface charge density σ at R , at a , and at b .
- Find the potential at the center, using infinity as the reference point.
- Now the outer surface is touched to a grounding wire, which lowers its potential to zero (same as at infinity). How do your answers to (a) and (b) change?

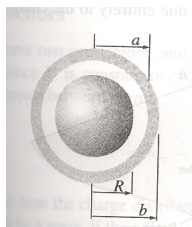
(a) $\sigma_R = \frac{q}{4\pi R^2}; \sigma_a = \frac{-q}{4\pi a^2}; \sigma_b = \frac{q}{4\pi b^2}.$

(b) $V(0) = -\int_{\infty}^0 \mathbf{E} \cdot d\mathbf{l} = -\int_{\infty}^b \left(\frac{1}{4\pi\epsilon_0} \frac{q}{r^2}\right) dr - \int_b^a (0) dr - \int_a^R \left(\frac{1}{4\pi\epsilon_0} \frac{q}{r^2}\right) dr - \int_R^0 (0) dr = \frac{1}{4\pi\epsilon_0} \left(\frac{q}{b} + \frac{q}{R} - \frac{q}{a}\right).$

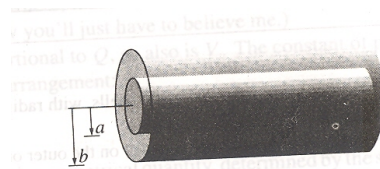
(c) $\sigma_b \rightarrow 0$ (the charge "drains off"); $V(0) = -\int_{\infty}^a (0) dr - \int_a^R \left(\frac{1}{4\pi\epsilon_0} \frac{q}{r^2}\right) dr - \int_R^0 (0) dr = \frac{1}{4\pi\epsilon_0} \left(\frac{q}{R} - \frac{q}{a}\right).$



(a) Problem 3



(b) Problem 6



(c) Problem 7

7. [G 2.39] Find the capacitance per unit length of two coaxial metal cylindrical tubes, of radii a and b (Fig.).

Say the charge on the inner cylinder is Q , for a length L . The field is given by Gauss's law:
 $\int \mathbf{E} \cdot d\mathbf{a} = E \cdot 2\pi s \cdot L = \frac{1}{\epsilon_0} Q_{\text{enc}} = \frac{1}{\epsilon_0} Q \Rightarrow E = \frac{Q}{2\pi\epsilon_0 L} \frac{1}{s} \hat{s}$. Potential difference between the cylinders is

$$V(b) - V(a) = -\int_a^b \mathbf{E} \cdot d\mathbf{l} = -\frac{Q}{2\pi\epsilon_0 L} \int_a^b \frac{1}{s} ds = -\frac{Q}{2\pi\epsilon_0 L} \ln\left(\frac{b}{a}\right).$$

As set up here, a is at the higher potential, so $V = V(a) - V(b) = \frac{Q}{2\pi\epsilon_0 L} \ln\left(\frac{b}{a}\right)$.

$$C = \frac{Q}{V} = \frac{2\pi\epsilon_0 L}{\ln\left(\frac{b}{a}\right)}, \text{ so capacitance per unit length is } \boxed{\frac{2\pi\epsilon_0}{\ln\left(\frac{b}{a}\right)}}.$$

8. [G 2.42] If the electric field in some region is given (in spherical coordinates) by the expression $\mathbf{E}(\mathbf{r}) = \frac{A\hat{\mathbf{r}} + B \sin \theta \cos \phi \hat{\phi}}{r}$, where A and B are constants, what is the charge density?

$$\begin{aligned}\rho &= \epsilon_0 \nabla \cdot \mathbf{E} = \epsilon_0 \left\{ \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{A}{r} \right) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \phi} \left(\frac{B \sin \theta \cos \phi}{r} \right) \right\} \\ &= \epsilon_0 \left[\frac{1}{r^2} A + \frac{1}{r \sin \theta} \frac{B \sin \theta}{r} (-\sin \phi) \right] = \boxed{\frac{\epsilon_0}{r^2} (A - B \sin \phi)}.\end{aligned}$$

9. A charge distribution on the surface of a sphere of radius R , is given by

$$\sigma(\theta, \phi) = \sigma_0 \sin \theta \sin \phi$$

Find the electric field and the potential on the z axis.

Let $\mathbf{r} = (0, 0, z)$ and $\mathbf{r}' = (R \sin \theta \cos \phi, R \sin \theta \sin \phi, R \cos \theta)$. Then $|\mathbf{r} - \mathbf{r}'| = (R^2 + z^2 - 2Rz \cos \theta)^{1/2}$.

$$\begin{aligned}E_x &= -\frac{\sigma_0 R^2}{4\pi\epsilon_0} \int_0^\pi \frac{\sin^3 \theta d\theta}{(R^2 + z^2 - 2Rz \cos \theta)^{3/2}} \underbrace{\int_0^{2\pi} \cos \phi \sin \phi d\phi}_{=0} = 0 \\ E_y &= -\frac{\sigma_0 R^2}{4\pi\epsilon_0} \int_0^\pi \frac{\sin^3 \theta d\theta}{(R^2 + z^2 - 2Rz \cos \theta)^{3/2}} \int_0^{2\pi} \sin^2 \phi d\phi = \begin{cases} -\frac{\sigma_0 R^3}{3\epsilon_0 z^3} & z > R \\ -\frac{\sigma_0}{3\epsilon_0} & z < R \end{cases} \\ E_z &= \frac{\sigma_0 R^2}{4\pi\epsilon_0} \int_0^\pi \frac{(z - R \cos \theta) \sin^2 \theta d\theta}{(R^2 + z^2 - 2Rz \cos \theta)^{3/2}} \underbrace{\int_0^{2\pi} \sin \phi d\phi}_{=0} = 0\end{aligned}$$

Now the θ integral in E_y is quiet messy, but has a nice compact expression.

10. Use Gauss's law to find electric field at any point due to the charge distribution

- (a) $\rho = \frac{\rho_0 a}{r}$ if $r < a$ and is zero otherwise;
 (b) $\rho = \rho_0(1 - r^2/a^2)$ if $r < a$ and is zero otherwise.

(a) Use a spherical Gaussian surface with centre at origin and radius r , then $\oint \mathbf{E} \cdot d\mathbf{S} = E \cdot 4\pi r^2$ and

$$q_{\text{enclosed}} = \rho_0 a \int_0^{\tilde{r}} \frac{1}{r} 4\pi r^2 dr = 4\pi \rho_0 a \frac{\tilde{r}^2}{2}$$

where $\tilde{r} = r$ if $r < a$, else $\tilde{r} = a$ if $r > a$. Then

$$\mathbf{E} = \begin{cases} \frac{\rho_0 a}{2\epsilon_0} \hat{\mathbf{r}} & r < a \\ \frac{\rho_0 a^3}{2\epsilon_0 r^2} \hat{\mathbf{r}} & r > a \end{cases}$$

(b) Simillary,

$$q_{\text{enclosed}} = \rho_0 \int_0^{\tilde{r}} \left(1 - \frac{r^2}{a^2}\right) 4\pi r^2 dr = 4\pi \rho_0 \tilde{r}^3 \left(\frac{1}{3} - \frac{1}{5} \frac{\tilde{r}^2}{a^2}\right)$$

Thus,

$$\mathbf{E} = \begin{cases} \frac{\rho_0 r}{3\epsilon_0} \left(1 - \frac{3}{5} \frac{r^2}{a^2}\right) \hat{\mathbf{r}} & r < a \\ \frac{2\rho_0 a^3}{15\epsilon_0 r^2} \hat{\mathbf{r}} & r > a \end{cases}$$