

CE 601: Numerical Methods

# Lecture 39

## Galerkin FEM for Laplace Equation-1

Course Coordinator:  
Dr. Suresh A. Kartha,  
Associate Professor,  
Department of Civil Engineering,  
IIT Guwahati.

## Galerkin FEM

Here, we will see how Galerkin FEM can be applied for 2-D cases.

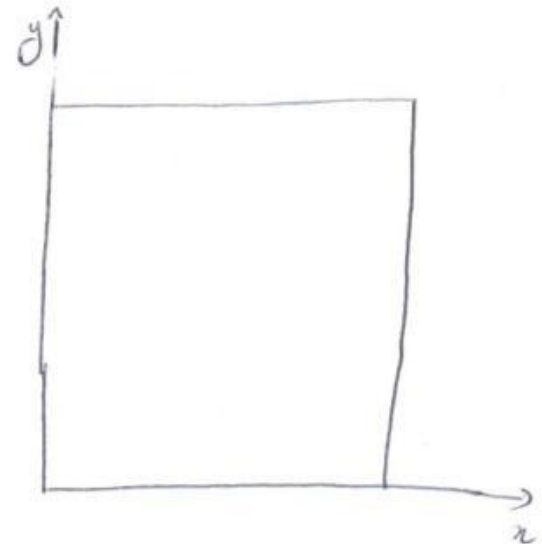
→ We will see Galerkin FEM to solve 2-D Laplace equation  
(or Poisson equation).

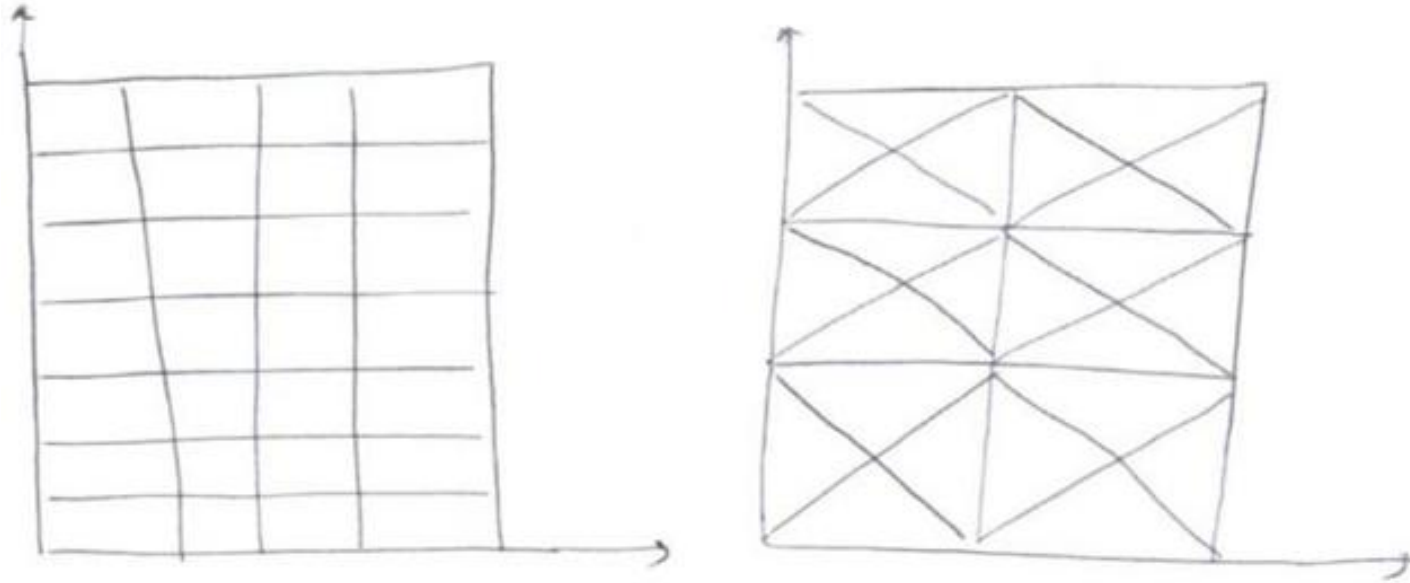
$$\left. \begin{array}{l} \frac{\partial^2 f}{\partial x^2} + \frac{\partial^2 f}{\partial y^2} = F(x, y) \\ \text{i.e.} \quad \quad \quad \text{or} \\ \frac{\partial^2 f}{\partial x^2} + \frac{\partial^2 f}{\partial y^2} = 0 \end{array} \right\} \text{with appropriate BCs.}$$

→ Let the two dimensional domain be rectangular in shape.

→ Figure shows the domain.

→ First thing is we need to discretise the domain into small small elements.



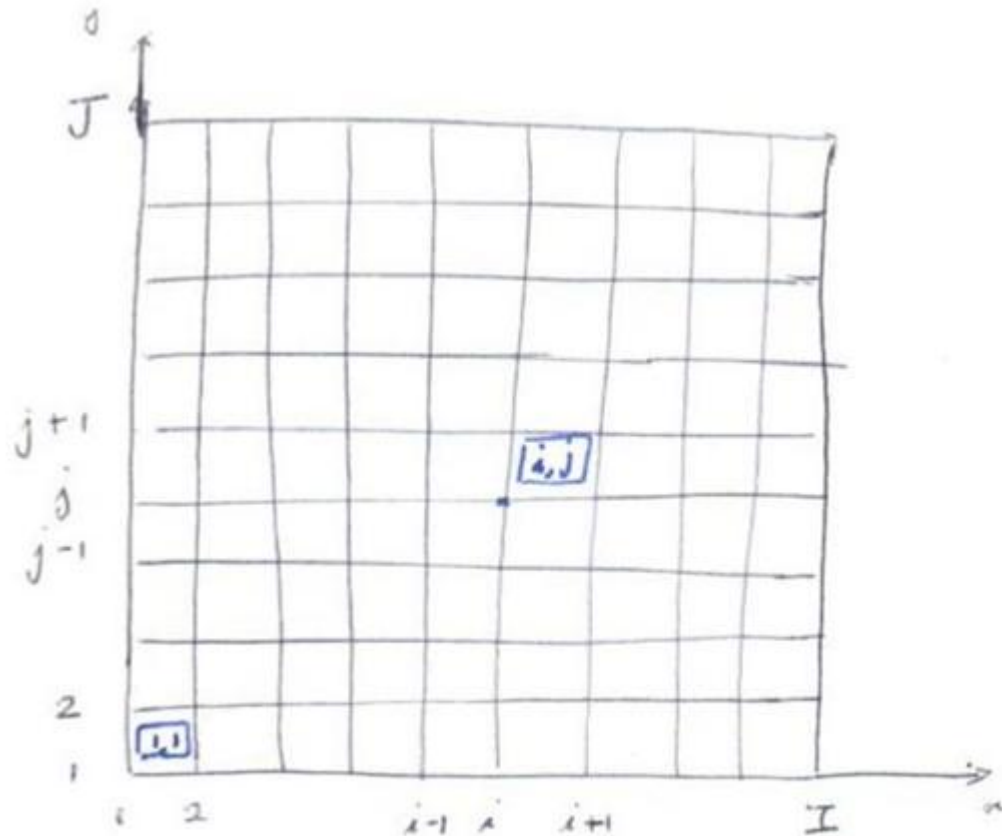


- The domain discretised using "rectangular elements", "triangular elements", "quadrilateral elements", etc.
- You have the freedom to choose any shape for the element.

→ In our case let us consider that the domain is discretised using rectangular elements.

∴ The  $x$  – axis is discretised  $i = 1, 2, 3, \dots, I$  nodes or discrete points

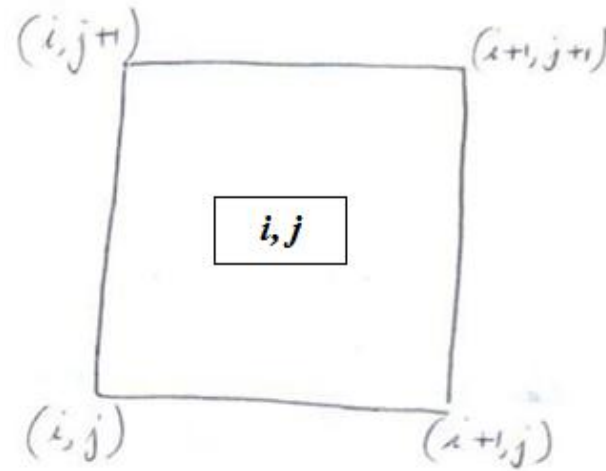
The  $y$  – axis is discretised  $j = 1, 2, 3, \dots, J$  nodes.



∴ For this rectangular domain there are a total of  $(I \times J)$  discrete nodes.

→ Any general node is given as  $(i, j)$  in the suffix.

→ Four nodes constitute an rectangular element. There are a total of  $(I - 1) \times (J - 1)$  elements. Any general element is given as  $[i, j]$  in the superfix.



→ As described for the one-dimensional case, there will be combination of approximating polynomial for 'f' in the entire domain

$$\text{i.e. } f(x, y) \approx \tilde{f}(x, y) = \sum_{i=1}^{I-1} \sum_{j=1}^{J-1} \tilde{f}^{[i,j]}(x, y)$$

(i.e. Sum of series of local interpolating polynomials)

→ For an element  $\boxed{i, j}$  → the approximation is  $\tilde{f}^{[i,j]}(x, y)$

→ As there are four nodes  $(i, j), (i+1, j), (i+1, j+1)$  and  $(i, j+1)$  associated with the element  $\boxed{i, j}$ , the approximation is given as:

$$\tilde{f}^{[i,j]}(x, y) = f_1 N_1(x, y) + f_2 N_2(x, y) + f_3 N_3(x, y) + f_4 N_4(x, y)$$

i.e. for the element  $\boxed{i, j}$ , the global node numbers are replaced by local node numbers 1, 2, 3, 4 and  $N_1, N_2, N_3, N_4$  are shape functions for the element  $\boxed{i, j}$ .

The shape functions are:

$$N_1(x, y) = \begin{cases} 1.0 & \text{at node 1} \\ 0 & \text{at other nodes} \end{cases}, N_2(x, y) = \begin{cases} 1.0 & \text{at node 2} \\ 0 & \text{at other nodes} \end{cases}$$

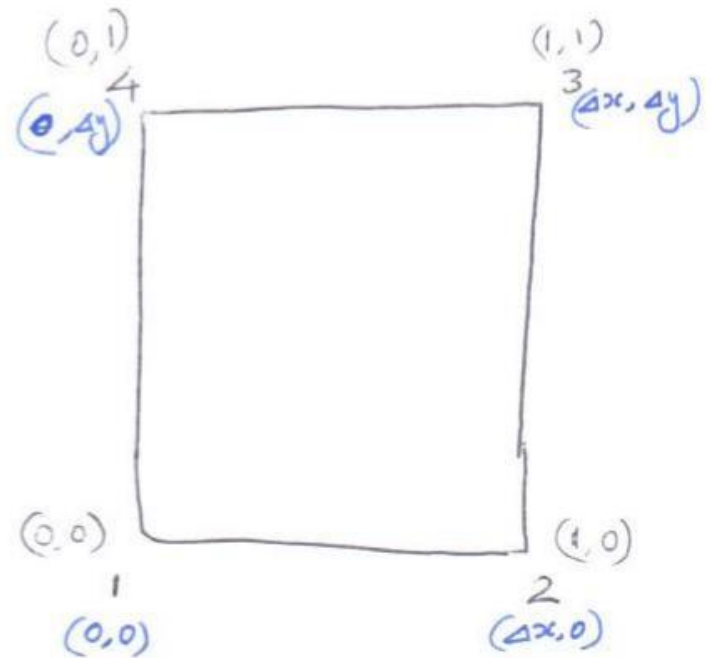
$$N_3(x, y) = \begin{cases} 1.0 & \text{at node 3} \\ 0 & \text{at other nodes} \end{cases}, N_4(x, y) = \begin{cases} 1.0 & \text{at node 4} \\ 0 & \text{at other nodes} \end{cases}$$

→ Let us assume the shape functions are linear

i.e.  $N(x, y) = a_0 + a_1x + a_2y + a_3xy$

→ As the shape functions for the element  $[i, j]$  have to be formalised, we use the normalised values of  $x$  and  $y$  within the element.

$$\text{i.e. } \bar{x} = \frac{x}{\Delta x} \text{ and } \bar{y} = \frac{y}{\Delta y}$$



→ The local coordinates at the node  $(i, j) \rightarrow (x, y) = (0, 0)$ ,  
 $(i + 1, j) \rightarrow (\Delta x, 0)$ ,  $(i + 1, j + 1) \rightarrow (\Delta x, \Delta y)$ , and  $(i, j + 1) \rightarrow (0, \Delta y)$ .

→ The normalised coordinates are therefore for the local element  
 $(0, 0) \rightarrow (0, 0)$ ,  $(\Delta x, 0) \rightarrow (1, 0)$

$(\Delta x, \Delta y) \rightarrow (1, 1)$  and  $(0, \Delta y) \rightarrow (0, 1)$

We develop shape functions using normalised coordinates  $(\bar{x}, \bar{y})$ .

$$N_1(\bar{x}, \bar{y}) = a_0 + a_1 \bar{x} + a_2 \bar{y} + a_3 \bar{x} \bar{y}$$

$$N_1(0, 0) = 1.0 = a_0 + 0 + 0 + 0; \therefore a_0 = 1.0$$

$$N_1(1, 0) = 0.0 = a_0 + a_1 + 0 + 0; \therefore a_1 = -1.0$$

$$N_1(0, 1) = 0.0 = a_0 + 0 + a_2 + 0; \therefore a_2 = -1.0$$

$$N_1(1, 1) = 0.0 = a_0 + a_1 + a_2 + a_3; \therefore a_3 = 1.0$$

$$\therefore N_1(\bar{x}, \bar{y}) = 1.0 - \bar{x} - \bar{y} + \bar{x} \bar{y}$$



Similarly,  $N_2(\bar{x}, \bar{y}) = a_0 + a_1\bar{x} + a_2\bar{y} + a_3\bar{x} \bar{y}$

We have  $N_2(0,0) = 0.0 = a_0; \therefore a_0 = 0.0$

$$N_2(1,0) = 1.0 = a_0 + a_1; \therefore a_1 = 1.0$$

$$N_2(0,1) = 0.0 = a_0 + a_2; \therefore a_2 = 0.0$$

$$N_2(1,1) = 0.0 = a_0 + a_1 + a_2 + a_3; \therefore a_3 = -1.0$$

i.e.  $N_2(\bar{x}, \bar{y}) = \bar{x} - \bar{x} \bar{y}$

Similarly,  $N_3(\bar{x}, \bar{y}) = \bar{x} \bar{y}$

and  $N_4(\bar{x}, \bar{y}) = \bar{y} - \bar{x} \bar{y}$

$$\begin{aligned} \therefore \tilde{f}^{[i,j]}(x, y) &= f_1 \times (1 - \bar{x} - \bar{y} + \bar{x} \bar{y}) + f_2 \times (\bar{x} - \bar{x} \bar{y}) + f_3 \times (\bar{x} \bar{y}) \\ &\quad + f_4 \times (\bar{y} - \bar{x} \bar{y}) \end{aligned}$$

This is the interpolating polynomial for the element  $\boxed{i, j}$ .

→ Define the Residual for the problem and evaluate weighted integral.

$$R(x, y) = \frac{\partial^2 f}{\partial x^2} + \frac{\partial^2 f}{\partial y^2} - F$$

$$\therefore I(f(x, y)) = \int_{x_1}^{x_l} \int_{y_1}^{y_l} W(x, y) \times \left( \frac{\partial^2 f}{\partial x^2} + \frac{\partial^2 f}{\partial y^2} - f \right) dy dx = 0$$

To do integration: see  $W \frac{\partial^2 f}{\partial x^2} = \frac{\partial}{\partial x} \left( W \frac{\partial f}{\partial x} \right) - \frac{\partial W}{\partial x} \frac{\partial f}{\partial x}$

$$\text{and } W \frac{\partial^2 f}{\partial y^2} = \frac{\partial}{\partial y} \left( W \frac{\partial f}{\partial y} \right) - \frac{\partial W}{\partial y} \frac{\partial f}{\partial y}$$

$$\therefore I = \int \int \left[ \frac{\partial}{\partial x} \left( W \frac{\partial f}{\partial x} \right) + \frac{\partial}{\partial y} \left( W \frac{\partial f}{\partial y} \right) - \frac{\partial W}{\partial x} \frac{\partial f}{\partial x} - \frac{\partial W}{\partial y} \frac{\partial f}{\partial y} - F \right] dx dy = 0$$

Using Stokes' theorem

$$\iint \frac{\partial}{\partial x} \left( W \frac{\partial f}{\partial x} \right) dx dy = \oint W \frac{\partial f}{\partial x} n_x ds$$

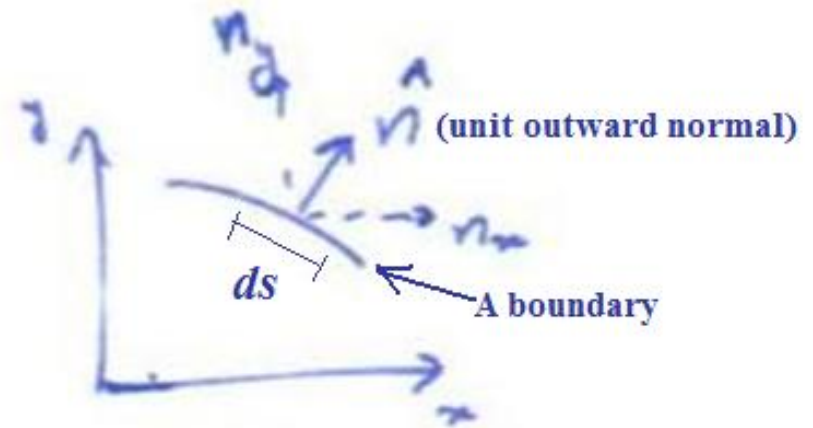
$$\iint \frac{\partial}{\partial y} \left( W \frac{\partial f}{\partial y} \right) dx dy = \oint W \frac{\partial f}{\partial y} n_y ds$$

where  $\hat{n} = n_x \hat{i} + n_y \hat{j}$

$$\therefore \iint \left[ \frac{\partial}{\partial x} \left( W \frac{\partial f}{\partial x} \right) + \frac{\partial}{\partial y} \left( W \frac{\partial f}{\partial y} \right) \right] dx dy = \oint \left( W \frac{\partial f}{\partial x} n_x + W \frac{\partial f}{\partial y} n_y \right) ds$$

Note that  $\hat{n} \cdot \nabla f = (n_x \hat{i} + n_y \hat{j}) \cdot \left( \frac{\partial f}{\partial x} \hat{i} + \frac{\partial f}{\partial y} \hat{j} \right)$

$$= n_x \frac{\partial f}{\partial x} + n_y \frac{\partial f}{\partial y} \equiv \text{flux of } f \text{ through boundary} = q_n$$



$$\begin{aligned}
\therefore I &= \iint \left[ \frac{\partial}{\partial x} \left( W \frac{\partial f}{\partial x} \right) + \frac{\partial}{\partial y} \left( W \frac{\partial f}{\partial y} \right) - \frac{\partial W}{\partial x} \frac{\partial f}{\partial x} - \frac{\partial W}{\partial y} \frac{\partial f}{\partial y} - F \right] dx dy = 0 \\
&= \oint_B W q_n ds - \iint \left( \frac{\partial W}{\partial x} \frac{\partial f}{\partial x} + \frac{\partial W}{\partial y} \frac{\partial f}{\partial y} + WF \right) dx dy = 0
\end{aligned}$$

The line integral describes the flux  $q_n$  normal to the outer boundary  $B$  of the solution domain.

→ For all interior elements that do not coincide

with the outer boundary, you have  $\oint_B W q_n ds = 0$

→ For Neumann BCs you have values for  $\oint_B W q_n ds$