

Kinematic Wave Routing

In the last class, we discussed on the analytical method of obtaining solutions to kinematic wave problem.

i.e.

We solved the governing equations

$$\frac{\partial Q}{\partial x} + \frac{\partial A}{\partial t} = 0$$
$$S_0 = S_f$$

→ We found the kinematic wave celerity.

$$c_k = \frac{dx}{dt} = \frac{dQ}{dA}$$

Today, we will briefly see the finite-difference numerical method approach to solve the kinematic wave equation.

We will briefly say the philosophy of kinematic wave ~~equation~~ finite difference method.

Finite-Difference Method to Solve Kinematic Waves Problem

- * FDM is an approximate method to solve pde or ODE.
- * The partial derivatives or derivatives are approximated by algebraic expressions.

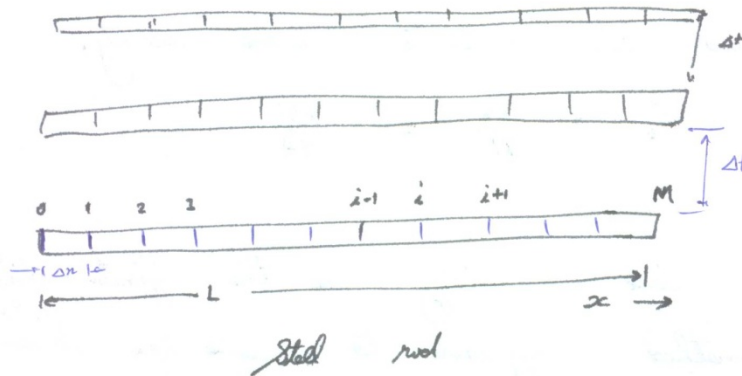
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* These these approximations are applied in pde (or ode) at any location in the domain to convert pde into algebraic equation.

* This algebraic equation is subsequently solved (by matrix methods, etc.)

To illustrate:

Consider a variable T (Temperature) that varies w.r.t space and time. For simplicity let us take a one dimensional steel rod



For some analysis we want to find the heat distribution in the steel rod.

As $T(x, t)$, $\rightarrow$ the heat problem may be function of $T, \frac{\partial T}{\partial x}, \frac{\partial T}{\partial t}, \frac{\partial^2 T}{\partial x^2}$, etc.

$\rightarrow$ Now the governing equation will involve the partial derivatives of T and T itself.

(3)

So to use finite difference method:

- (i) First we will discretize the spatial domain into finite small domains of length Δx
- (ii) We will assign node numbers at the interface $i = 1, 2, 3, \dots, M$
- (iii) Similarly we will try to evaluate the temperature at various time steps. That is the entire derived duration is broken into small time steps of Δt duration
- (iv) Now the temperature T at any spatial node i and time step s ~~will~~ can be represented

as:

$$T_i^{(s)}$$

- (v) The spatial partial derivative

$$\left. \frac{\partial T}{\partial x} \right|_i^{(s)} \approx \frac{T_{i+1}^{(s)} - T_{i-1}^{(s)}}{2\Delta x} \quad \text{Centered Difference}$$

$$\approx \frac{T_{i+1}^{(s)} - T_i^{(s)}}{\Delta x} \quad \text{Forward difference}$$

$$\approx \frac{T_i^{(s)} - T_{i-1}^{(s)}}{\Delta x} \quad \text{Backward difference}$$

- (vi) Similarly the temporal derivative

$$\left. \frac{\partial T}{\partial t} \right|_i^{(s)} \approx \frac{T_i^{(s+1)} - T_i^{(s)}}{\Delta t} \quad \text{Forward Time}$$

$$\approx \frac{T_i^{(s)} - T_i^{(s-1)}}{\Delta t} \quad \text{Backward Time.}$$

4

$$\frac{\partial^2 T}{\partial x^2} \bigg|_i^{(A)} = \frac{T_{i+1}^{(A)} - 2T_i^{(A)} + T_{i-1}^{(A)}}{\Delta x^2} \quad (\text{Central Difference}).$$

So in the governing pde, if we approximate the time derivative by forward time approximation $\rightarrow$ we generally get explicit expressions.

* If backward time difference is used, then we get implicit algebraic approximation for the pde.

For the Kinematic Wave

The continuity equation is

$$\frac{\partial Q}{\partial x} + \frac{\partial A}{\partial t} = q$$

$$S_0 = S_f$$

ie. We can write:

$$\frac{\partial Q}{\partial x} + \alpha \beta Q^{p-1} \frac{\partial Q}{\partial t} = q \rightarrow (1)$$

Now this is the governing ~~pt equation~~ pde for kinematic wave.

On solving this we will get $Q(x, t)$ for different values of x and time t .

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We are now approximating the partial derivative by linear scheme.

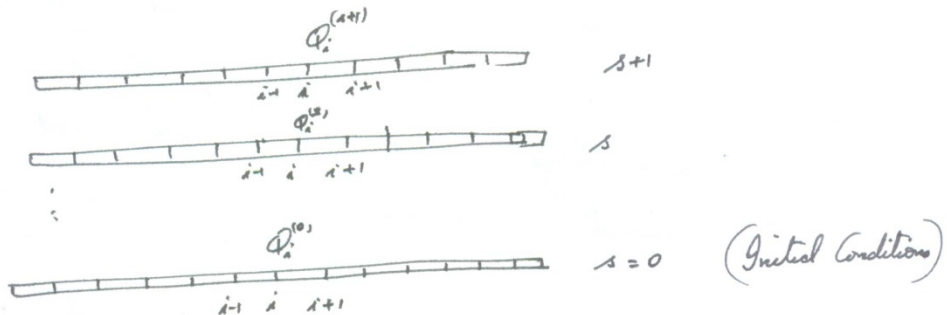
$$\frac{\partial \varphi}{\partial x} \Big|_i^{(s)} \approx \frac{\varphi_{i+1}^{(s)} - \varphi_{i-1}^{(s)}}{2 \Delta x}$$

or We will use forward difference:

$$\frac{\partial \varphi}{\partial x} \Big|_i^{(s)} = \frac{\varphi_{i+1}^{(s)} - \varphi_i^{(s)}}{\Delta x}$$

$$\frac{\partial \varphi}{\partial t} \Big|_i^{(s)} = \frac{\varphi_i^{(s+1)} - \varphi_i^{(s)}}{\Delta t}$$

→ Substitute these approximation in pde.



$$\alpha \frac{\partial \varphi}{\partial x} \Big|_i^{(s)} + \alpha \beta (\varphi_i^{(s)})^{\beta-1} \frac{\partial \varphi}{\partial t} \Big|_i^{(s)} = q_i^{(s)}$$

$$\frac{\varphi_{i+1}^{(s)} - \varphi_i^{(s)}}{\Delta x} + \alpha \beta (\varphi_i^{(s)})^{\beta-1} \frac{\varphi_i^{(s+1)} - \varphi_i^{(s)}}{\Delta t} = q_i^{(s)}$$

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∴ $Q_i^{(s+1)}$ = Unknown quantity

$$Q_i^{(s+1)} = Q_i^{(s)} + \frac{q_i^{(s)} \Delta t}{\alpha \beta (Q_i^{(s)})^{\beta-1}} - \frac{\Delta t}{\Delta x} \frac{1}{\alpha \beta (Q_i^{(s)})^{\beta-1}} [Q_{i+1}^{(s)} - Q_i^{(s)}]$$

Similarly if you take

$$\left. \frac{\partial Q}{\partial x} \right|_i^{(s+1/2)} + \alpha \beta (Q_i^{(s+1/2)})^{\beta-1} \left. \frac{\partial Q}{\partial t} \right|_i^{(s+1/2)} = q_i^{(s+1/2)}$$

The centered difference comes into picture.

$$Q_i^{(s+1/2)} = \frac{Q_i^{(s)} + Q_{i+1}^{(s)}}{2}$$

Muskingum - Cunge Method

You have seen Muskingum method for lumped flow routing. With same approach and improvement Muskingum - Cunge method can be used to route kinematic wave.

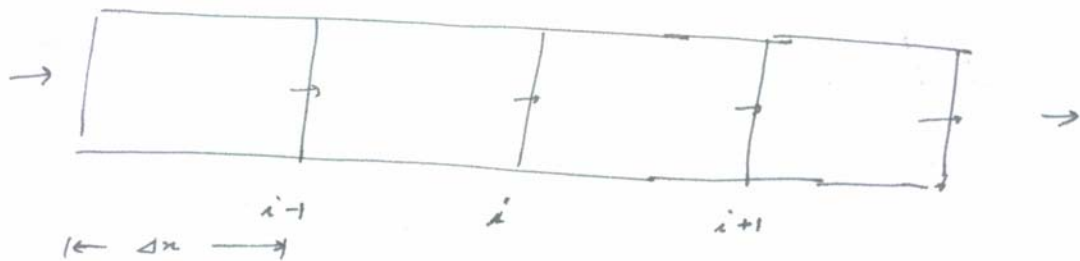
Recall

$$Q_{j+1} = C_1 I_{j+1} + C_2 I_j + C_3 Q_j$$

Here the suffix was to express time interval.

(7)

In Muskingum - lump method, what we will do is both spatial and temporal variations are incorporated.



→ The inflow hydrograph is given at $i = 0$.

$$\therefore P_{i+1}^{(s+1)} = C_1 P_i^{(s+1)} + C_2 P_i^{(s)} + C_3 P_{i+1}^{(s)}$$

$C_1, C_2, C_3 \rightarrow$ coefficients.

$K \rightarrow$ storage constant.

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DYNAMIC WAVE ROUTING

As you have seen many civilisations developed along river banks.

There are more urbanisations along river banks.

To ~~implement~~ ^{implement} routing we require more accurate predictions by routing.

$\therefore$ Dynamic wave routing is required.

$\rightarrow$ If backwater effects are negligible we were able to use kinematic wave approximation.

The dynamic wave equation:

$$\frac{\partial \Phi}{\partial t} + \frac{\partial}{\partial x} \left(\beta \frac{\Phi^2}{A} \right) + g A \left(\frac{\partial y}{\partial x} - S_0 + S_f + S_e \right) - \beta q v_u + W_f B = 0$$

(If all quantities are essential and then incorporated).

When the flow is uniform or $S_0 = S_f$ all other quantities negligible, you have.

i.e. Unique value of Φ w.r.t stage.

$\rightarrow$ However for dynamic condition, not unique.

