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Lecture 26

Differential Analysis of Flow (Contd..)

In the last class, we derived the partial differential equation for *conservation of mass* principle.



Fig. 1: Reference Coordinate axes

Recall that equation:

$$\frac{\partial \rho}{\partial t} + \frac{\partial (\rho u)}{\partial x} + \frac{\partial (\rho v)}{\partial y} + \frac{\partial (\rho w)}{\partial z} = 0 \qquad \rightarrow (1)$$
$$\vec{v} = u\hat{i} + v\hat{j} + w\hat{k}$$

where the velocity vector is or

$$\vec{v} = v_1 \hat{e}_1 + v_2 \hat{e}_2 + v_3 \hat{e}_3$$

Looking into equation 1, it is obvious to you that it can be expressed as

$$\frac{\partial \rho}{\partial t} + \nabla .(\rho \vec{v}) = 0 \qquad \rightarrow (2)$$

➢ By this time, you should be knowing, how to write equation 2 in index notation. $\frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x_i} (\rho v_i) = 0$

- > If the fluid is incompressible, then $\frac{\partial \rho}{\partial t} = 0$.
- > Therefore, the conservation of mass equation becomes

 $\nabla .(\rho \vec{v}) = 0$ i.e. $\rho(\nabla . \vec{v}) = 0$ $\nabla . \vec{v} = 0 \longrightarrow (3)$

That is, for incompressible flow, $\nabla . \vec{v} = 0$

If for any case, you need to use cylindrical polar coordinates, the principle will still be valid.



Fig.2: Definition sketch for cylindrical coordinate system (Source: Fluid Mechanics by F.M. White)

The velocity components are:

Axial velocity : vz

Radial velocity : vr

Circumferential velocity : v_{θ}

Like in Cartesian coordinates $(x_1, x_2, x_3 \text{ or } x, y, z)$ in cylindrical polar coordinates any fluid property will be function of (r, θ, z, t) .

How are you going to write the conservation of mass equation in cylindrical polar coordinates ?

Note the relation between cylindrical polar coordinates and Cartesian coordinates :

$$r = \sqrt{x^2 + y^2}$$
$$\theta = \tan^{-1}\left(\frac{y}{x}\right)$$
$$z = z$$

> The divergence of any vector \vec{B} in cylindrical polar coordinates is:

$$\vec{B} = B_r \hat{r} + B_\theta \hat{\theta} + B_z \hat{z}$$

$$\nabla .\vec{B} = \frac{1}{r} \frac{\partial (rB_r)}{\partial r} + \frac{1}{r} \frac{\partial (B_\theta)}{\partial \theta} + \frac{\partial (B_z)}{\partial z}$$
Therefore, the continuity equation will be :
$$\frac{\partial \rho}{\partial t} + \frac{1}{r} \frac{\partial (r\rho v_r)}{\partial r} + \frac{1}{r} \frac{\partial (\rho v_\theta)}{\partial \theta} + \frac{\partial (v_z)}{\partial z}$$

Conservation of Linear Momentum

The principle of conservation of linear momentum can also be differentially analyzed for fluids.

Here again, we will take the same rectangular prismoidal element that was discussed for conservation of mass principle.



- > This differential element has six rectangular faces.
- ➤ As per the shape we are assuming 1-D input and output faces.

Since the volume is fixed in space and stationary, the momentum principle using RTT for this elementary volume will be:

$$\frac{D(m\vec{v})}{Dt} = \frac{\partial}{\partial t} [\iiint_{cv} \vec{v} \rho dU] + \iint_{cs} \vec{v} \rho(\vec{v}.\hat{n}) dA$$

i.e. $\Sigma \vec{F} = \frac{\partial}{\partial t} [\iiint_{cv} \vec{v} \rho dU] + \sum_{outlet} [\rho |\vec{v}| A\vec{v}] - \sum_{inlet} \rho |\vec{v}| A\vec{v}$

Since there is 1-dimensional input and output, assumed.

- > The surface integral term is Net momentum outflux.
- > The volumetric integral:

$$\frac{\partial}{\partial t} [\iiint_{cv} \vec{v} \rho dU] \approx \frac{\partial (\rho \vec{v})}{\partial t} \Delta x \Delta y \Delta z$$

(for the rectangular elemental volume)

> Let the momentum influx on the left side will be = $\rho v \vec{v} \Delta x \Delta z$ And the momentum outflux will be = $[\rho v \vec{v} + \frac{\partial (\rho v \vec{v})}{\partial y} \Delta y] \Delta x \Delta z$

Therefore, the net momentum outflux in the y-direction will be outflux - influx

$$= \left[\rho v \vec{v} + \frac{\partial(\rho v v)}{\partial y} \Delta y\right] \Delta x \Delta z - \rho v \vec{v} \Delta x \Delta z$$
$$= \frac{\partial(\rho v \vec{v})}{\partial y} \Delta x \Delta y \Delta z$$

→

Similarly, in x-direction = $\frac{\partial(\rho u \vec{v})}{\partial x} \Delta x \Delta y \Delta z$ And, in z- direction = $\frac{\partial(\rho w \vec{v})}{\partial z} \Delta x \Delta y \Delta z$

Therefore, the RTT becomes:

$$\begin{split} \Sigma \vec{F} &= \frac{\partial(\rho \vec{v})}{\partial t} \Delta x \Delta y \Delta z + \frac{\partial(\rho u \vec{v})}{\partial x} \Delta x \Delta y \Delta z + \frac{\partial(\rho v \vec{v})}{\partial y} \Delta x \Delta y \Delta z + \frac{\partial(\rho w \vec{v})}{\partial z} \Delta x \Delta y \Delta z \\ \Sigma \vec{F} &= [\frac{\partial(\rho \vec{v})}{\partial t} + \frac{\partial(\rho u \vec{v})}{\partial x} + \frac{\partial(\rho v \vec{v})}{\partial y} + \frac{\partial(\rho w \vec{v})}{\partial z}] \Delta x \Delta y \Delta z \\ \Sigma \vec{F} &= [\vec{v} \frac{\partial\rho}{\partial t} + \vec{v} (\frac{\partial(\rho u)}{\partial x} + \frac{\partial(\rho v)}{\partial y} + \frac{\partial(\rho w)}{\partial z}) + \rho \frac{\partial \vec{v}}{\partial t} + \rho u \frac{\partial \vec{v}}{\partial x} + \rho v \frac{\partial \vec{v}}{\partial y} + \rho w \frac{\partial \vec{v}}{\partial z}] \Delta x \Delta y \Delta z \\ \Sigma \vec{F} &= [\vec{v} \{\frac{\partial\rho}{\partial t} + \nabla . (\rho \vec{v})\} + \rho \{\frac{\partial \vec{v}}{\partial t} + u \frac{\partial \vec{v}}{\partial x} + v \frac{\partial \vec{v}}{\partial y} + w \frac{\partial \vec{v}}{\partial z}\}] \Delta x \Delta y \Delta z \\ Note \quad \frac{\partial\rho}{\partial t} + \nabla . (\rho \vec{v}) = 0 \\ \text{i.e.} \qquad \Sigma \vec{F} &= \rho \Delta x \Delta y \Delta z [\frac{\partial \vec{v}}{\partial t} + (\vec{v} . \nabla) \vec{v}] \\ \text{Recall acceleration, } \vec{a} &= \frac{D \vec{v}}{Dt} = \frac{\partial \vec{v}}{\partial t} + (\vec{v} . \nabla) \vec{v} \text{ (studied earlier in particle kinematics)} \\ \Sigma \vec{F} &= e^{D \vec{v}} \Delta x \Delta y \Delta z$$

$$\sum \vec{F} = \rho \frac{Dv}{Dt} \Delta x \Delta y \Delta z$$

- > As discussed earlier there can be:
 - ✤ Body forces
 - Surface forces
- > The body forces can be:
 - 1. Gravitational
 - 2. Magnetic
 - 3. Electric potential, etc.
- ▶ For our usual fluid like water, etc. the magnetic and electric potential are neglected.

Quiz:

- 1. Write in index notation, the conservation of mass and conservation of momentum equation.
- 2. Write the conservation of mass equation in cylindrical polar coordinates.