

LINEAR MOMENTUM IN DIFFERENTIAL APPROACH
(contd...)

Yesterday, we started discussing about linear momentum principle through differential approach.

We have seen for an elemental volume Δxyz

$$\vec{F} = \rho \Delta xyz \left[\frac{\partial \vec{v}}{\partial t} + (\vec{v} \cdot \nabla) \vec{v} \right]$$

$$\text{or } \vec{F} = \rho \frac{d\vec{v}}{dt} \Delta xyz$$

→ There will be - Body forces and surface forces that contributes to \vec{F}

→ For this situation, let us consider the body force is only due to gravity

$$\therefore \vec{f}_{grav} = \rho \vec{g} \Delta xyz$$

$$\text{where } \vec{g} = \rho \hat{i} + \rho \hat{j} - g \hat{k}$$

→ The surface force will be due to pressure and viscous stresses.

→ The surface forces can be related with respect to the stresses on the sides of the rectangular element considered here.

These stresses are sum of hydrostatic pressure and viscous stress.

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Viscous stress can be represented as $\bar{\tau}$ or
in short form τ_{ij} . They are tensorial
quantities.

e.g. Velocity is a vector

$$\vec{v} = u\hat{i} + v\hat{j} + w\hat{k} \rightarrow v_i$$

i.e. It has three components.

In other words, we require three components
to define a vector.
Similarly a ^{second rank} tensorial quantity requires nine components
to be described.

$$\bar{\tau} = \begin{pmatrix} \tau_{xx} & \tau_{xy} & \tau_{xz} \\ \tau_{yx} & \tau_{yy} & \tau_{yz} \\ \tau_{zx} & \tau_{zy} & \tau_{zz} \end{pmatrix} \rightarrow \tau_{ij}$$

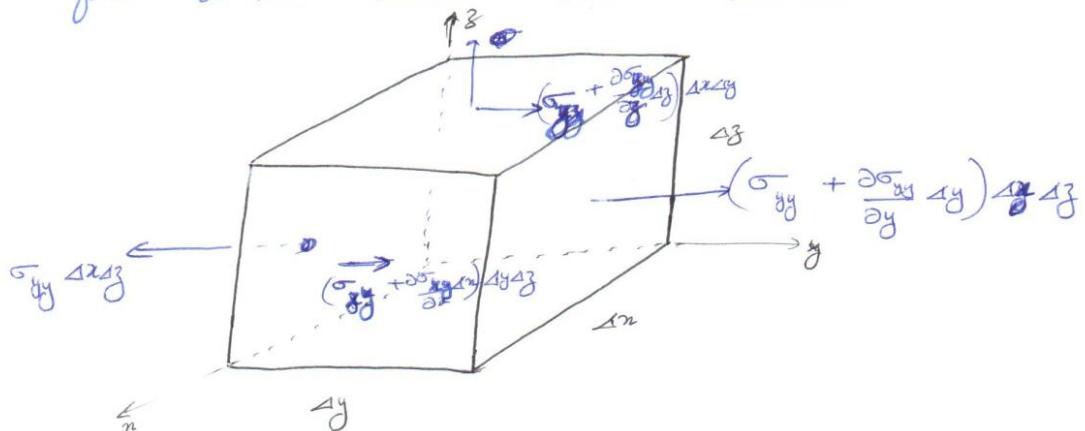
Similarly, strain rates E_{ij} are also second rank
tensors.

The stress, therefore, is a second rank tensor

$$\bar{\sigma} \rightarrow \sigma_{ij} = \begin{pmatrix} -p + \tau_{xx} & \tau_{yx} & \tau_{zx} \\ \tau_{xy} & -p + \tau_{yy} & \tau_{zy} \\ \tau_{xz} & \tau_{yz} & -p + \tau_{zz} \end{pmatrix}$$

(3)

As discussed earlier for pressure, here also, it is the gradients or differences that cause a net ~~for~~ surface force on the elemental control volume.



Consider ~~again~~ the surface force, right now, only in y -direction you have normal forces $\left\{ \left(\sigma_{yy} + \frac{\partial \sigma_{yy}}{\partial y} \Delta y \right) \Delta x \Delta z - \sigma_{yy} \Delta x \Delta z \right\}$

$$= \frac{\partial \sigma_{yy}}{\partial y} \Delta x \Delta y \Delta z$$

You have ^{also} tangential or shear forces in y -direction

$$\left(\sigma_{xy} + \frac{\partial \sigma_{xy}}{\partial x} \Delta x \right) \Delta y \Delta z - \sigma_{xy} \Delta y \Delta z$$

and $\left(\sigma_{zy} + \frac{\partial \sigma_{zy}}{\partial z} \Delta z \right) \Delta x \Delta y - \sigma_{zy} \Delta x \Delta y$

\therefore The Net ^{surface} force in x -direction

$$dF_{y, \text{surface}} = \left[\frac{\partial (\sigma_{xy})}{\partial x} + \frac{\partial \sigma_{yy}}{\partial y} + \frac{\partial \sigma_{zy}}{\partial z} \right] \Delta x \Delta y \Delta z$$

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$$\frac{dF_{y, \text{surface}}}{du} = -\frac{\partial p}{\partial y} + \frac{\partial (\tau_{xy})}{\partial x} + \frac{\partial (\tau_{yy})}{\partial y} + \frac{\partial (\tau_{yz})}{\partial z}$$

$$16) \quad \frac{dF_{x, \text{surface}}}{du} = -\frac{\partial p}{\partial x} + \frac{\partial (\tau_{xx})}{\partial x} + \frac{\partial (\tau_{xy})}{\partial y} + \frac{\partial (\tau_{xz})}{\partial z}$$

$$\frac{dF_{z, \text{surface}}}{du} = -\frac{\partial p}{\partial z} + \frac{\partial (\tau_{xz})}{\partial x} + \frac{\partial (\tau_{yz})}{\partial y} + \frac{\partial (\tau_{zz})}{\partial z}$$

$$\hat{i} \frac{dF_{x, \text{surface}}}{du} + \hat{j} \frac{dF_{y, \text{surface}}}{du} + \hat{k} \frac{dF_{z, \text{surface}}}{du} = \vec{dF}_{\text{surface}}$$

$$\frac{\vec{dF}_{\text{surface}}}{du} = -\nabla p + \frac{\vec{dF}_{\text{viscous}}}{du}$$

Now

$$\frac{\vec{dF}_{\text{viscous}}}{du} = i \left[\frac{\partial e}{\partial x} + \frac{\partial e}{\partial y} + \frac{\partial e}{\partial z} \right] + j \left[\frac{\partial e}{\partial x} + \frac{\partial e}{\partial y} + \frac{\partial e}{\partial z} \right] + k \left[\frac{\partial e}{\partial x} + \frac{\partial e}{\partial y} + \frac{\partial e}{\partial z} \right]$$

$$= \bar{\epsilon} \cdot \bar{\tau}$$

, please note $\bar{\tau}$ is a
second rank tensor

$$\bar{\tau} = \begin{pmatrix} \tau_{xx} & \tau_{xy} & \tau_{xz} \\ \tau_{yx} & \tau_{yy} & \tau_{yz} \\ \tau_{zx} & \tau_{zy} & \tau_{zz} \end{pmatrix}$$

(5)

$\bar{\tau}$ \rightarrow Various stress tensors acting on the elemental fluid volume $\Delta x \Delta y \Delta z$.

\therefore We have now from the relation:

$$\vec{F} = \rho \frac{d\vec{v}}{dt} \Delta x \Delta y \Delta z$$

$$\text{i.e. } \vec{dF}_{\text{grav}} + \vec{dF}_{\text{surface}} = \rho \frac{d\vec{v}}{dt} \Delta x \Delta y \Delta z$$

$$\text{i.e. } \rho \vec{g} \Delta x \Delta y \Delta z + [(-\nabla p) \Delta x \Delta y \Delta z + \nabla \cdot \bar{\tau} \Delta x \Delta y \Delta z] = \rho \frac{d\vec{v}}{dt} \Delta x \Delta y \Delta z$$

As the volume is arbitrary, the equation becomes
expression at any mathematical point.

$$\boxed{\rho \vec{g} - \nabla p + \nabla \cdot \bar{\tau} = \rho \frac{d\vec{v}}{dt}} \rightarrow (1)$$

⑥

$$\begin{aligned} \oint g_x - \frac{\partial p}{\partial x} + \frac{\partial z_x}{\partial x} e + \frac{\partial z_y}{\partial x} e + \frac{\partial z_z}{\partial x} e \\ = \oint \left(\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} + w \frac{\partial u}{\partial z} \right) \end{aligned}$$

$$\begin{aligned} \oint g_y - \frac{\partial p}{\partial y} + \frac{\partial z_x}{\partial y} e + \frac{\partial z_y}{\partial y} e + \frac{\partial z_z}{\partial y} e \\ = \oint \left[\frac{\partial v}{\partial t} + u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} + w \frac{\partial v}{\partial z} \right] \end{aligned}$$

$$\begin{aligned} \oint g_z - \frac{\partial p}{\partial z} + \frac{\partial z_x}{\partial z} e + \frac{\partial z_y}{\partial z} e + \frac{\partial z_z}{\partial z} e \\ = \oint \left[\frac{\partial w}{\partial t} + u \frac{\partial w}{\partial x} + v \frac{\partial w}{\partial y} + w \frac{\partial w}{\partial z} \right] \end{aligned}$$