

[Griffiths] 7.8, 7.11, 7.16, 7.22, 7.26, 7.38, 7.41, 7.52

[Griffiths] 8.2, 8.4, 9.9, 9.10, 9.12, 9.13, 9.15,

Answers

G7.8 (a) Flux through the square loop = $\frac{\mu_0 I a}{2\pi} \ln\left(\frac{s+a}{s}\right)$.

(b) emf = $\frac{\mu_0 I a^2 v}{2\pi s(s+a)}$. Current flows counterclockwise in the figure.

(c) emf = 0.

G7.11 Emf in loop $\mathcal{E} = Blv$. Current $I = Blv/R$. Upward force on loop = $IlB = B^2 l^2 v/R$. Then the equation of motion is

$$m \frac{dv}{dt} = mg - \frac{B^2 l^2 v}{R}.$$

Upon solving,

$$v = \frac{mgR}{B^2 l^2} (1 - e^{-\alpha t})$$

where $\alpha = \frac{B^2 l^2}{mR}$.

G7.16 (a) Suppose current in the inside wire is in \hat{k} direction. The magnetic field is circumferential. Thus electric field is in \hat{k} direction.

(b) $\mathbf{E}(s, \phi, z) = \frac{\mu_0 I_0 \omega}{2\pi} \sin(\omega t) \ln(a/s) \hat{k}$.

G7.22 Flux through single turn = $\mu_0 n I \pi R^2$. In one unit length n such turns. The total flux = $\mu_0 n^2 I \pi R^2$. Self Inductance = $\mu_0 n^2 \pi R^2$

G7.26 (a) Inductance per unit length = $\mu_0 n^2 \pi R^2$. Energy per unit length $W = \frac{1}{2} L I^2 = \mu_0 n^2 \pi R^2 I^2 / 2$.

(b) $W = \frac{1}{2} \oint (\mathbf{A} \cdot \mathbf{I}) dl$. Now, $A = (\mu_0 n I R / 2) \hat{\phi}$ on the surface. Then $W = \frac{1}{2} ((\mu_0 n I R / 2) I (2\pi R))$

(c) $W = \frac{1}{2\mu_0} \int B^2 d\tau$. Now, $B = \mu_0 n I$ inside and zero outside. Take a volume with one unit length along the axis. Then $W = \frac{1}{2\mu_0} (\mu_0 n I)^2 \pi R^2$.

G7.38 The figure shows a hemispherical surface. The current through the surface is $I = \int J \cdot da = \sigma \int E \cdot da$. Here σ is the conductivity of the material. The surface can be converted a closed Gaussian surface by closing it in the conductor, where E is zero. Thus $I = \sigma Q_{enc} / \epsilon_0$. Where $Q_{enc} = \int \sigma_e da$. However, (From eq 3.77) $\sigma_e = 3\epsilon_0 E_0 \cos \theta$. Here $E_0 = V_0/d$. Thus $I = 3\sigma \pi V_0 a^2 / d$.

G7.41 The potential:

$$\begin{aligned} V_i(s, \phi) &= \sum_{k=1}^{\infty} c_k s^k \sin(k\phi) & s < a \\ V_o(s, \phi) &= \sum_{k=1}^{\infty} b_k s^{-k} \sin(k\phi) & s > a \end{aligned}$$

Cosines will not contribute since the potential function is odd on the surface.
At $s = a$, $V_i = V_o = V_0\phi/2\pi$. Thus

$$\begin{aligned}c_k &= \frac{V_0}{2\pi^2 a^k} \int_{-\pi}^{\pi} \phi \sin(k\phi) d\phi = \frac{V_0}{2\pi^2 a^k} \left[-\frac{2\pi}{k} (-1)^k \right] \\b_k &= \frac{V_0 a^k}{2\pi^2} \int_{-\pi}^{\pi} \phi \sin(k\phi) d\phi = \frac{V_0 a^k}{2\pi^2} \left[-\frac{2\pi}{k} (-1)^k \right]\end{aligned}$$

then

$$\begin{aligned}V_i(s, \phi) &= -\frac{V_0}{\pi} \sum_{k=1}^{\infty} \frac{1}{k} \left(-\frac{s}{a}\right)^k \sin(k\phi) = \frac{V_0}{\pi} \tan^{-1} \left(\frac{s \sin \phi}{a + s \cos \phi} \right) \\V_o(s, \phi) &= -\frac{V_0}{\pi} \sum_{k=1}^{\infty} \frac{1}{k} \left(-\frac{a}{s}\right)^k \sin(k\phi) = \frac{V_0}{\pi} \tan^{-1} \left(\frac{a \sin \phi}{s + a \cos \phi} \right)\end{aligned}$$

G7.52 A point on the lower loop: $\mathbf{r}_1 = (b \cos \phi_1, b \sin \phi_1, 0)$ and a point on upper loop: $\mathbf{r}_2 = (a \cos \phi_2, a \sin \phi_2, z)$. Then

$$\begin{aligned}r^2 &= (\mathbf{r}_1 - \mathbf{r}_2)^2 \\&= a^2 + b^2 + z^2 - 2ab \cos(\phi_1 - \phi_2) \\&= \frac{ab}{\beta} [1 - 2\beta \cos(\phi_1 - \phi_2)]\end{aligned}$$

$d\mathbf{l}_1 = b d\phi_1 \hat{\phi}_1$; $d\mathbf{l}_2 = a d\phi_2 \hat{\phi}_2$. Then $d\mathbf{l}_1 \cdot d\mathbf{l}_2 = ab d\phi_1 d\phi_2 \hat{\phi}_1 \cdot \hat{\phi}_2 = ab d\phi_1 d\phi_2 \cos(\phi_1 - \phi_2)$.

$$M = \frac{\mu_0}{4\pi} \oint \oint \frac{d\mathbf{l}_1 \cdot d\mathbf{l}_2}{r} = \frac{\mu_0}{4\pi} \frac{ab}{\sqrt{ab/\beta}} \int \int \frac{\cos(\phi_1 - \phi_2)}{\sqrt{1 - 2\beta \cos(\phi_1 - \phi_2)}} d\phi_1 d\phi_2$$

With some simplification

$$M = \frac{\mu_0}{4\pi} \sqrt{ab\beta} \int_0^{2\pi} \frac{\cos u}{\sqrt{1 - 2\beta \cos u}} du$$

If a is small compared to b and z , then $\beta \ll 1$. Use Binomial theorem to obtain results given in the book.

G8.2 Discussed in class.

G8.4 Let the positive charge be at $(0, 0, a)$ and the negative charge at $(0, 0, -a)$. The two regions are $z > 0$ and $z < 0$. To find the force on the positive charge we will consider the region $z > 0$. The Maxwells stress tensor is given by $T_{ij} = \epsilon_0(E_i E_j - \delta_{ij} E^2/2)$, since this is electrostatic case with no magnetic field. Thus on $z = 0$ surface, $E(r, \pi/2, \phi) = (0, 0, E_z)$ Where $E_z = \frac{2qa}{4\pi\epsilon_0(r^2+a^2)^{3/2}}$

$$T = \begin{pmatrix} -\epsilon_0 E_z^2/2 & 0 & 0 \\ 0 & -\epsilon_0 E_z^2/2 & 0 \\ 0 & 0 & \epsilon_0 E_z^2/2 \end{pmatrix}$$

Now the force on the upper region is given by $F = \oint_s T \cdot da$. Since $da = -rdrd\phi\hat{k}$, the force is

$$\begin{aligned} F &= - \int T_{zz} r dr d\phi \\ &= -\frac{1}{2} \left(\frac{2qa}{4\pi\epsilon_0} \right)^2 \epsilon_0 \int \frac{rdrd\phi}{(r^2 + a^2)^3} \\ &= -\frac{q}{4\pi\epsilon_0} \frac{1}{4a^2} \end{aligned}$$

G9.9 (a)

$$\begin{aligned} \mathbf{E}(x, y, z, t) &= \hat{\mathbf{z}} E_0 \cos\left(\frac{\omega}{c}x + \omega t\right) \\ \mathbf{B}(x, y, z, t) &= \hat{\mathbf{y}} \frac{E_0}{c} \cos\left(\frac{\omega}{c}x + \omega t\right) \end{aligned}$$

(b) $\mathbf{k} = \frac{\omega}{\sqrt{3}c}(1, 1, 1)$ and $\mathbf{n} = \frac{1}{\sqrt{2}}(\hat{\mathbf{x}} - \hat{\mathbf{z}})$, Thus $\mathbf{k} \times \mathbf{n} = \frac{1}{\sqrt{6}}(-\hat{\mathbf{x}} + 2\hat{\mathbf{y}} - \hat{\mathbf{z}})$.

$$\begin{aligned} \mathbf{E}(x, y, z, t) &= E_0 \cos\left(\frac{\omega}{\sqrt{3}c}(x + y + z) - \omega t\right) \left(\frac{\hat{\mathbf{x}} - \hat{\mathbf{z}}}{\sqrt{2}}\right) \\ \mathbf{B}(x, y, z, t) &= \frac{E_0}{\sqrt{6}c} \cos\left(\frac{\omega}{\sqrt{3}c}(x + y + z) - \omega t\right) (-\hat{\mathbf{x}} + 2\hat{\mathbf{y}} - \hat{\mathbf{z}}) \end{aligned}$$

G9.10 Pressure $P = I/c = 4.3 \times 10^{-6} \text{N/m}^2$.

G9.12 $\mathbf{E}(x, y, z, t) = \hat{\mathbf{x}} E_0 \cos(\frac{\omega}{c}z - \omega t)$ and $\mathbf{B}(x, y, z, t) = \hat{\mathbf{y}} \frac{E_0}{c} \cos(\frac{\omega}{c}z - \omega t)$. Thus,

$$T_{xx} = \epsilon_0 \left(E_x E_x - \frac{1}{2} E^2 \right) + \frac{1}{\mu_0} \left(-\frac{1}{2} B^2 \right) = \frac{1}{2} \left(\epsilon_0 E^2 - \frac{1}{\mu_0} B^2 \right) = 0$$

Similarly for $T_{yy} = 0$. And

$$T_{zz} = \epsilon_0 \left(-\frac{1}{2} E^2 \right) + \frac{1}{\mu_0} \left(-\frac{1}{2} B^2 \right) = -\frac{1}{2} \left(\epsilon_0 E^2 + \frac{1}{\mu_0} B^2 \right) = -u$$

G9.13

G9.15 If the given equation is true for all x , the derivative of the equation is also true at all x . Then

$$\begin{array}{ll} A + B &= C \quad \text{put } x=0 \\ aA + bB &= cC \quad \text{differentiate and put } x=0 \\ a^2A + b^2B &= c^2C \quad \text{differentiate twice and put } x=0 \end{array}$$

For non-trivial solution for A, B and C , the determinant of coefficients must be zero. This quickly gives $a = b$. In addition all A, B and C must be nonzero gives $b = c$.