

1. Prove the following identities:

(a) $[A, B + C] = [A, B] + [A, C]$.

(b) $[A, BC] = [A, B]C + B[A, C]$.

(c) $[A, [B, C]] + [B, [C, A]] + [C, [A, B]] = 0$.

(d) $[f(\hat{X}), \hat{P}] = i\hbar \frac{df}{dX}$, where f is operator function of \hat{X} .

2. For any operator A , show that $(A + A^\dagger)$, $i(A - A^\dagger)$ and AA^\dagger are hermitian operators.

3. Prove:

$$e^{+A} B e^{-A} = B + [A, B] + \frac{1}{2!} [A, [A, B]] + \frac{1}{3!} [A, [A, [A, B]]] + \dots$$

and if A and B commute with their commutator then

$$e^A e^B = e^{A+B+\frac{1}{2}[A,B]}$$

4. Let \mathcal{H} be a Hilbert space. Prove Schwarz inequality, that is, for any two vectors, $f, g \in \mathcal{H}$, show that

$$|\langle f, g \rangle|^2 \leq \langle f, f \rangle \langle g, g \rangle.$$

[Hint: Consider $\langle f + \lambda g, f + \lambda g \rangle \geq 0$. Now find λ such that the lhs is minimum.]

5. Let a time dependent observable be represented by a hermitian operator $\hat{\Omega}(t)$. If the system is in state $\Psi(t)$, show that

$$\frac{d}{dt} \langle \hat{\Omega} \rangle = \frac{i}{\hbar} \langle [\hat{H}, \hat{\Omega}] \rangle + \left\langle \frac{\partial \hat{\Omega}}{\partial t} \right\rangle$$

where \hat{H} is the hamiltonian operator. [Hint: If $f(t)$ and $g(t)$ are any two states, then prove that

$$\frac{d}{dt} \langle f(t), g(t) \rangle = \left\langle \frac{d}{dt} f(t), g(t) \right\rangle + \left\langle f(t), \frac{d}{dt} g(t) \right\rangle$$

using the properties of inner product and

$$\frac{d}{dt} \langle f(t), g(t) \rangle = \lim_{\Delta t \rightarrow 0} \frac{\langle f(t + \Delta t), g(t + \Delta t) \rangle - \langle f(t), g(t) \rangle}{\Delta t}$$

6. If the hamiltonian operator of a system is given by $\hat{H} = \hat{P}^2/2m + V(\hat{X})$, then prove the Ehrenfest's theorem:

$$\begin{aligned} \frac{d}{dt} \langle \hat{X} \rangle &= \frac{1}{m} \langle \hat{P} \rangle \\ \frac{d}{dt} \langle \hat{P} \rangle &= \langle -V'(\hat{X}) \rangle. \end{aligned}$$

[Hint: Use the commutation relations $[\hat{X}, \hat{P}] = i\hbar$ and the properties of operators.]

7. Consider a quantum system consisting of a particle in a conservative force field. The energy spectrum is $\{E_1, E_2, \dots\}$ with corresponding normalized stationary states $\{\phi_1, \phi_2, \dots\}$. Let x_0 be an eigenvalue of the position operator with eigenvector ξ_{x_0} . Let $\alpha_1 = \langle \xi_{x_0}, \phi_1 \rangle$ and $\alpha_2 = \langle \xi_{x_0}, \phi_2 \rangle$. Let $\Psi(t)$ denote the state of the system at time t . Express, in terms of α_1, α_2 and eigenenergies, the answers to the following questions:

- (a) If $\Psi(0) = \phi_1$, what is the probability density of finding the particle at x_0 at time t ?
 (b) If $\Psi(0) = \phi_2$, what is the probability density of finding the particle at x_0 at time t ?
 (c) If $\Psi(0) = (\phi_1 + \phi_2)/\sqrt{2}$, what is the probability density of finding the particle at x_0 at time t ? What is the maximum probability density? And minimum?

8. Here is an example of time dependent Hamiltonian: An electron in an oscillating electric field is described by a Hamiltonian operator

$$\hat{H} = \frac{\hat{P}^2}{2m} - (eE_0 \cos \omega t) x$$

where E_0 is the amplitude of the electric field. Calculate $d\langle \hat{X} \rangle / dt$, and $d\langle \hat{P} \rangle / dt$.

9. The potential energy of a harmonic oscillator is given by $V(x) = m\omega^2 x^2 / 2$. Assuming that $\langle \hat{X} \rangle = \langle \hat{P} \rangle = 0$, find the lower limit to the expectation value of the Hamiltonian operator. [Hint: Use the uncertainty principle.]
10. The first excited state of the harmonic oscillator is given by

$$\psi_1(x) = \left(\frac{2\alpha}{\sqrt{\pi}} \right)^{1/2} (\alpha x) e^{-\alpha^2 x^2 / 2}$$

Find ΔX and ΔP and check uncertainty principle. $(\alpha = \frac{m\omega}{\hbar})$ $(\alpha^2 = \frac{m\omega}{\hbar})$

Tutorial 3.

Q1. (a) $[A, B+C] = A(B+C) - (B+C)A = (AB-BA) + (AC-CA) = [A, B] + [A, C].$

(b) $[A, BC] = ABC - BCA = ABC - BAC + BAC - BCA = [A, B]C + B[A, C].$

(c) Use brute force.

(d) $[f(\hat{x}), \hat{p}]g(x) = [f(\hat{x})\hat{p} - \hat{p}f(\hat{x})]g(x)$
 $= f(x) \cdot (-i\hbar \frac{d}{dx})g(x) - (-i\hbar \frac{d}{dx})f(x)g(x)$
 $= f(x)\hat{p}g(x) - f(x)\hat{p}g(x) - (-i\hbar \frac{df}{dx})g(x)$
 $= i\hbar \frac{df}{dx}(x)g(x)$

$\Rightarrow [f(\hat{x}), \hat{p}] = i\hbar \frac{df}{dx}$

Q2. $(A+A^\dagger)^\dagger = A^\dagger + (A^\dagger)^\dagger = A^\dagger + A.$

$[i(A-A^\dagger)]^\dagger = (A^\dagger - (A^\dagger)^\dagger) i^\dagger = (-i)(A^\dagger - A) = i(A-A^\dagger)$

$(AA^\dagger)^\dagger = A^\dagger(A^\dagger)^\dagger = A^\dagger A.$

Q3. Step 1: Prove that $\frac{d}{d\lambda} e^{\lambda A} = A e^{\lambda A}$ (remember since A is an operator this identity is not obvious.)
 using $\frac{d}{d\lambda} e^{\lambda A} = \lim_{\Delta\lambda \rightarrow 0} \frac{e^{(\lambda+\Delta\lambda)A} - e^{\lambda A}}{\Delta\lambda}$

Step 2: Consider

$f(\lambda) = e^{\lambda A} B e^{-\lambda A}$

then $f(\lambda) = f(0) + \lambda f'(0) + \frac{\lambda^2}{2!} f''(0) + \dots$ ①

Now $f(0) = B$, $f'(0) = \left. \frac{d}{d\lambda} (e^{\lambda A} B e^{-\lambda A}) \right|_{\lambda=0}$
 $= \left. (A e^{\lambda A} B e^{-\lambda A} + e^{\lambda A} B (-A e^{-\lambda A})) \right|_{\lambda=0}$
 $= e^{\lambda A} [A, B] e^{-\lambda A} \Big|_{\lambda=0}$
 $= [A, B]$

and $f''(0) = [A, [A, B]]$ etc.

Now, substitute in ① and put $\lambda=1$.

Q4. Now $\langle f+\lambda g, f+\lambda g \rangle \geq 0$ true for all λ

$$\Rightarrow \underbrace{|f|^2 + |\lambda|^2 |g|^2 + \lambda \langle f, g \rangle + \lambda^* \langle g, f \rangle}_{\text{lhs}} \geq 0$$

Now lhs will be minimum when

$$\lambda = - \frac{\langle g, f \rangle}{|g|^2}$$

(while finding min, remember λ is complex No.)

Then

$$\Rightarrow |f|^2 + \frac{|\langle f, g \rangle|^2}{|g|^4} |g|^2 - \frac{|\langle f, g \rangle|^2}{|g|^2} - \frac{|\langle f, g \rangle|^2}{|g|^2} \geq 0$$

$$\Rightarrow \underline{|f|^2 |g|^2 \geq |\langle f, g \rangle|^2}$$

Q5. Step 1. Now

$$\begin{aligned} \frac{d}{dt} \langle f(t), g(t) \rangle &= \lim_{\Delta t \rightarrow 0} \frac{\langle f(t+\Delta t), g(t+\Delta t) \rangle - \langle f(t), g(t) \rangle}{\Delta t} \\ &= \lim_{\Delta t \rightarrow 0} \left\{ \frac{\langle f(t+\Delta t), g(t+\Delta t) \rangle - \langle f(t), g(t+\Delta t) \rangle}{\Delta t} \right. \\ &\quad \left. + \frac{\langle f(t), g(t+\Delta t) \rangle - \langle f(t), g(t) \rangle}{\Delta t} \right\} \\ &= \lim_{\Delta t \rightarrow 0} \left[\left\langle \frac{f(t+\Delta t) - f(t)}{\Delta t}, g(t+\Delta t) \right\rangle + \left\langle f(t), \frac{g(t+\Delta t) - g(t)}{\Delta t} \right\rangle \right] \\ &= \left\langle \frac{df}{dt}, g \right\rangle + \left\langle f, \frac{dg}{dt} \right\rangle. \end{aligned}$$

Step 2.

$$\begin{aligned} \frac{d}{dt} \langle \hat{Q}(t) \rangle &= \frac{d}{dt} \langle \hat{\Psi}(t), \hat{Q}(t) \hat{\Psi}(t) \rangle \\ &= \left\langle \frac{d\hat{\Psi}}{dt}, \hat{Q}(t) \hat{\Psi}(t) \right\rangle + \left\langle \hat{\Psi}(t), \left(\frac{d\hat{Q}}{dt} \right) \hat{\Psi}(t) \right\rangle \\ &\quad + \left\langle \hat{\Psi}, \hat{Q} \left(\frac{d\hat{\Psi}}{dt} \right) \right\rangle \end{aligned}$$

By Schrödinger eq. $i\hbar \frac{d\hat{\Psi}}{dt} = \hat{H} \hat{\Psi}$

Then

$$\begin{aligned}\frac{d}{dt} \langle \hat{Q}(t) \rangle &= \left\langle \frac{1}{i\hbar} \hat{H} \psi, \hat{Q} \psi \right\rangle + \left\langle \bar{\psi}, \hat{Q} \left(\frac{1}{i\hbar} \hat{H} \psi \right) \right\rangle \\ &\quad + \left\langle \bar{\psi}, \left(\frac{\partial \hat{Q}(t)}{\partial t} \right) \psi \right\rangle \\ &= i\hbar^{-1} \langle \hat{H} \hat{Q} - \hat{Q} \hat{H} \rangle + \left\langle \frac{\partial \hat{Q}(t)}{\partial t} \right\rangle \\ &= i\hbar^{-1} \langle [\hat{H}, \hat{Q}(t)] \rangle + \left\langle \frac{\partial \hat{Q}(t)}{\partial t} \right\rangle.\end{aligned}$$

Q6. Using the result of (5)

$$\begin{aligned}[\hat{H}, \hat{X}] &= \frac{1}{2m} [P^2, \hat{X}] = \frac{1}{2m} \{ P [\hat{P}, \hat{X}] + [\hat{P}, \hat{X}] P \} \\ &= \frac{1}{2m} (2\hat{P} (i\hbar)) = -i\hbar \frac{\hat{P}}{m}\end{aligned}$$

$$[\hat{H}, \hat{P}] = [V(\hat{X}), \hat{P}] = i\hbar \frac{dV}{dx}$$

Then

$$\frac{d}{dt} \langle \hat{X} \rangle = \frac{i}{\hbar} \langle [\hat{H}, \hat{X}] \rangle = \frac{i}{\hbar} (-i\hbar) \frac{\langle \hat{P} \rangle}{m} = \frac{\langle \hat{P} \rangle}{m}$$

$$\begin{aligned}\frac{d}{dt} \langle \hat{P} \rangle &= \frac{i}{\hbar} \langle [\hat{H}, \hat{P}] \rangle = \frac{i}{\hbar} (i\hbar) \langle \frac{dV}{dx} \rangle \\ &= \langle -V'(x) \rangle.\end{aligned}$$

Q7. (a) if $\Psi(0) = \phi_1$ then $\Psi(t) = e^{-iE_1 t/\hbar} \phi_1$

$$P_{\hat{X}}(x_0, t) = |\langle \xi_{x_0}, \Psi(t) \rangle|^2 = |\langle \xi_{x_0}, \phi_1 \rangle|^2 = \alpha_1^2$$

(b) Similarly

$$P_{\hat{X}}(x_0, t) = \alpha_2^2$$

(c) $\Psi(0) = \frac{1}{\sqrt{2}} (\phi_1 + \phi_2)$

$$\text{Then } \Psi(t) = \frac{1}{\sqrt{2}} \left[\phi_1 e^{-iE_1 t/\hbar} + \phi_2 e^{-iE_2 t/\hbar} \right]$$

Thus

$$\begin{aligned}
 P_x(x_0, t) &= \left| \langle \xi_{x_0}, \frac{1}{\sqrt{2}} \phi_1 e^{-iE_1 t/\hbar} \rangle + \langle \xi_{x_0}, \frac{1}{\sqrt{2}} \phi_2 e^{-iE_2 t/\hbar} \rangle \right|^2 \\
 &= \frac{1}{2} \alpha_1^2 + \frac{1}{2} \alpha_2^2 + 2 \operatorname{Re} \left[\frac{1}{2} \langle \xi_{x_0}, \phi_1 \rangle^* \langle \xi_{x_0}, \phi_2 \rangle e^{-i(E_2 - E_1)t/\hbar} \right. \\
 &\quad \left. + \frac{1}{2} \langle \xi_{x_0}, \phi_1 \rangle \langle \xi_{x_0}, \phi_2 \rangle^* e^{-i(E_1 - E_2)t/\hbar} \right] \\
 &= \frac{1}{2} \alpha_1^2 + \frac{1}{2} \alpha_2^2 + \frac{1}{2} \cdot 2 \cdot \operatorname{Re} (\alpha_1 \alpha_2^* e^{-i(E_1 - E_2)t/\hbar})
 \end{aligned}$$

Max when $e^{-i(E_1 - E_2)t/\hbar} = +1$

$$= \frac{1}{2} (\alpha_1 + \alpha_2)^2$$

Min when $e^{-i(E_1 - E_2)t/\hbar} = -1$

$$= \frac{1}{2} (\alpha_1 - \alpha_2)^2$$

Q8.

$$\frac{d\langle \hat{x} \rangle}{dt} = \frac{\langle \hat{p} \rangle}{m}$$

$$\frac{d\langle \hat{p} \rangle}{dt} = +e E_0 \cos \omega t$$

Q9.

$$\langle \hat{H} \rangle = \frac{\langle \hat{p}^2 \rangle}{2m} + \frac{1}{2} m \omega^2 \langle \hat{x}^2 \rangle$$

Thus $\sigma_H = \sqrt{\langle \hat{H}^2 \rangle - \langle \hat{H} \rangle^2}$

10. Given $\psi_1(x) = \left(\frac{2\alpha}{\sqrt{\pi}}\right)^{1/2} e^{-\alpha x} e^{-\alpha^2 x^2/2}$

(a) $\langle x \rangle = 0$ $\langle p \rangle = 0$

(b) $\langle x^2 \rangle = \left(\frac{2\alpha}{\sqrt{\pi}}\right) \alpha \int_{-\infty}^{\infty} (\psi_1(x))^2 e^{-\alpha^2 x^2} x^2 dx$
 $= \frac{3}{2} \alpha^2$

(c) $\langle p^2 \rangle = \frac{3\alpha^2 \hbar^2}{2}$

(d) $\Delta x = \sqrt{\frac{3}{2}} \frac{1}{\alpha}$ $\Delta p = \sqrt{\frac{3}{2}} \alpha \hbar$

$\Delta x \Delta p = \frac{3}{2} \hbar > \frac{1}{2} \hbar$